

# ITER Baseline Scenario – Q=10 operation

by  
**F. Turco**  
with  
*...many others*

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# What is a scenario

..., for our purpose today:

- **plasma goals**
  - power output, gain, duration, ...
- **boundary conditions**
  - heating mix, plasma formation, evolution, ...
- **constraints set by machine parameters**
  - flux consumption, tolerable li, coil currents, divertor loads, ...

# What is the ITER Baseline Scenario (affectionately, the IBS)



## IBS Mission:

- $P_{\text{fus}}=500 \text{ MW} \rightarrow \beta_{\text{T,therm}}=2.55\%$   
 $\rightarrow \beta_{\text{N}} \sim 1.8$  (low  $\beta_{\text{N}}$ )
- $Q=10 = P_{\text{fus}} / (P_{\text{transp}} + P_{\text{loss}} - P_{\alpha}) \rightarrow \beta_{\text{T}}\tau_{\text{E}}$  or  $G = \frac{\beta_{\text{N}}H}{q_{95}^2} \sim 0.4$  (proxy)  
 $\rightarrow$  Minimise input power  $\rightarrow$  **Need high confinement**  $\rightarrow$  High  $I_p$
- At full field  $B_{\text{T}}=5.3 \text{ T}$ ,  $I_p=15 \text{ MA}$   
 $\rightarrow q_{95} \sim 3$  (lower current alternatives require higher  $\beta_{\text{N}}$  or  $H_{98}$ )

## Constraints for demonstration discharges:

- ITER shape (affects MHD stability, pedestal)
- Zero injected torque (moderate rotation)
- H-mode on  $I_p$  flatter (heated ramps are possible)

# What we discuss today

1. Standard IBS pulse design
2. MHD stability and disruptions
  - Cause of the instabilities
  - Solution and new scenario
3. Confinement trends
4. Lower current and heated access options

# What we discuss today

## 1. Standard IBS pulse design

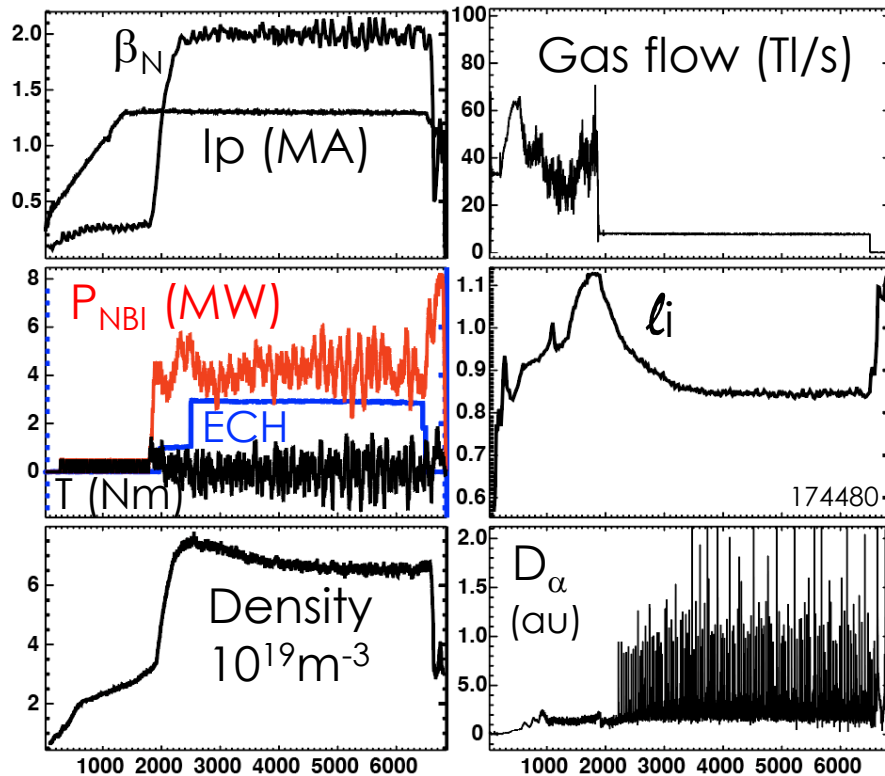
## 2. MHD stability and disruptions

- Cause of the instabilities
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## 3. Confinement trends

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# The IBS is Designed with an Ohmic Ramp-up, Zero Injected Torque Throughout and Low $\beta_N$

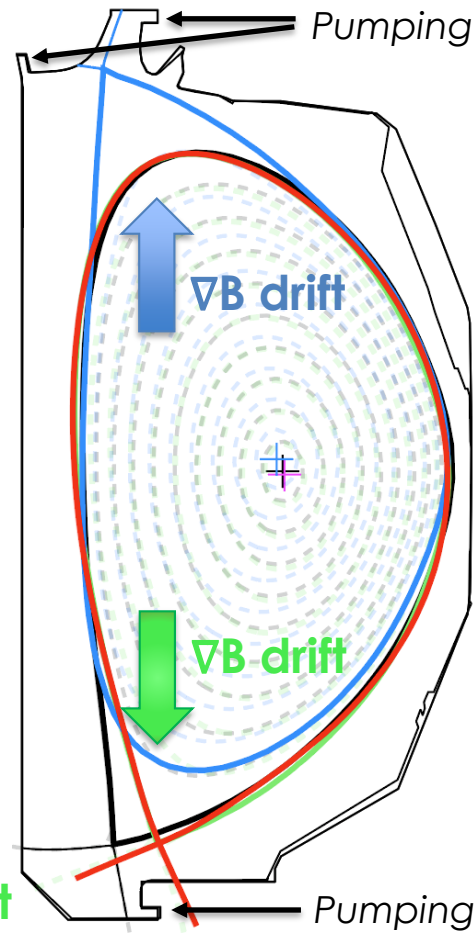


— ITER (x0.2778)

— DIII-D pumped Std Bt

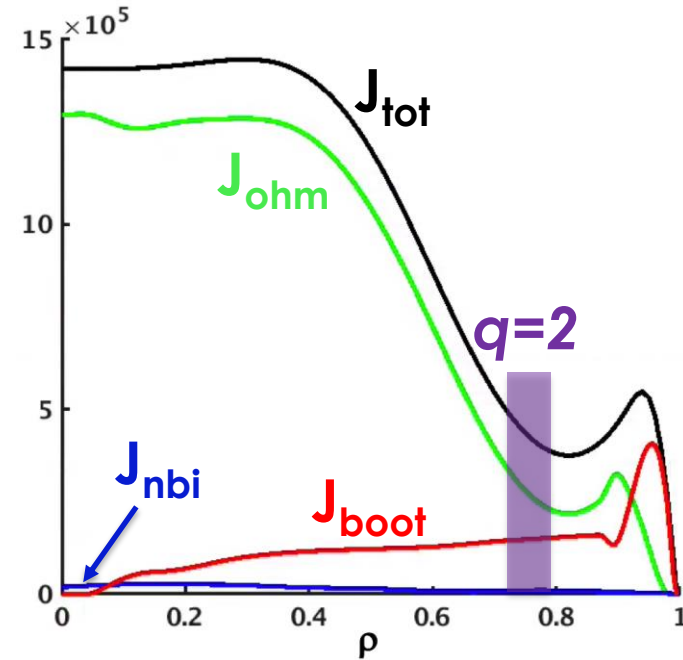
— DIII-D pumped Rev Bt

— DIII-D unpumped Std Bt



# Characteristics of the Equilibrium

- The plasma current is mostly **inductive**
  - High  $I_p$ , low  $\beta_N$  for high gain
  - Active current tailoring is not possible
- $J_{boot}$  dominates the pedestal
- $q_0 \sim 1$  (sawteeth),  $q_{95} \sim 3 \rightarrow q=2$  at  $\rho \sim 0.8 \rightarrow$  Strong correlation between pedestal and core



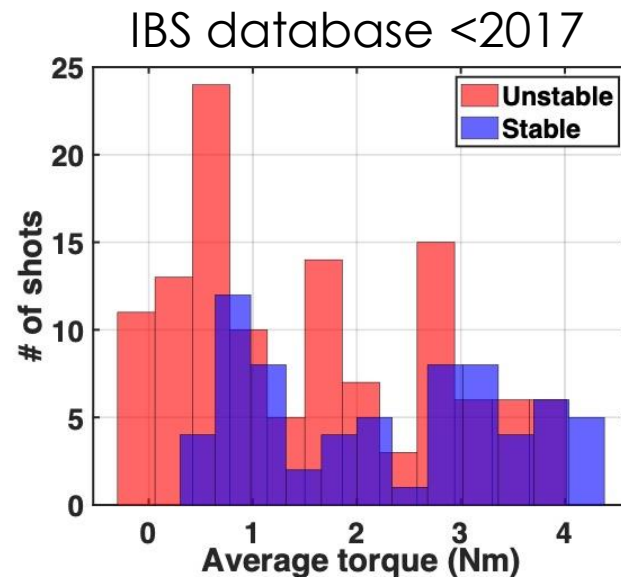
# For 15 Years, the IBS in DIII-D Was Often Terminated by a Fast Growing, Disrupting 2/1 Tearing Mode

## Before 2017

- Zero stable IBS at  $T < 0.5$  Nm
- Over 65% of full co-torque shots disrupt

There are no instabilities in the ramp-up (ohmic)  
There can be some in the pre-programmed ramp-down

**Today I will be discussing the instabilities that occur on the  $\beta_N$  flat-top (burn phase)**



**Keep in mind! We are interested in the ONSET, not the growth** (the modes collapse the pressure and disrupt in 10-200 ms)



# What we discuss today

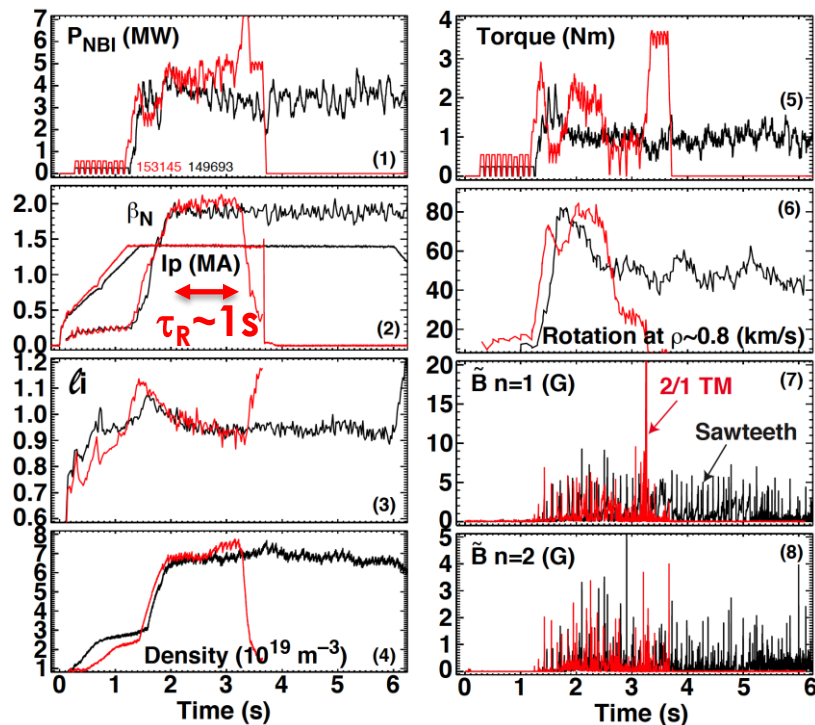
1. Standard IBS pulse design
- 2. MHD stability and disruptions**
  - **Cause of the instabilities**
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# This 2/1 TM Is Not Your Garden Variety Pressure-Driven Instability

## Characteristics of these modes:

- They occur at fixed  $\beta_N$ , pressure shape
- After several  $\tau_{ES}$  (pressure equilibrated)
- After tens of sawteeth, hundreds of ELMs (not seeded)
- At all torque, rotation and  $\nabla\Omega$  values
- Lower  $\beta_N$  is not better
- While the current profile  $J$  is evolving

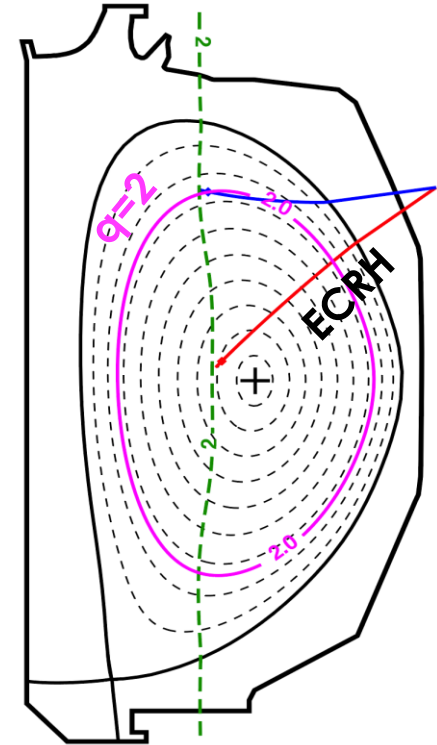
*Studying the details of the TM drives led to a solution for robustly stable plasmas*



# Direct ECCD Stabilization is Not a Viable Option for This Scenario

- For this low- $\beta_N$  scenario active stabilization of TMs has not been successful (and not for lack of trying)
- The TMs are likely *not* neoclassical in nature  $\rightarrow$  replacing  $j_{\text{boot}}$  does not eliminate the drive ( $\Delta'$ )
- $q=2$  is near the edge  $\rightarrow$  ECCD at  $\rho \sim 0.75$  is very detrimental to the performance (shown later)

Different methods to achieve stability have to be investigated



# Which Physics Mechanisms Can Affect the ONSET of Tearing modes?

- Pressure  $p \rightarrow$  it evolves on the  $\tau_E$  time scale ( $\sim 100$  ms in DIII-D)
- Current profile  $J \rightarrow$  it evolves on the  $\tau_R$  time scale ( $\sim 1$  s in these plasmas)
- Rotation  $\Omega$  and  $\nabla\Omega \rightarrow$  it evolves on the  $\tau_E$  time scale
- Mode coupling  $\rightarrow$  perturbed field from  $n>1$  modes may resonate with 2/1 rational surface
- Seeding  $\rightarrow$  it assumes the mode is classically stable, but "noise" produces SMALL islands at 2/1 surface (sawteeth, ELMs)

*The typical tearing time is  $\tau_{\text{tear}} \sim 5$  ms  $\ll \tau_E, \tau_R$   
If  $p, J, \Omega$ , "seeds" are right, it MUST tear in  $\leq 5$  ms*

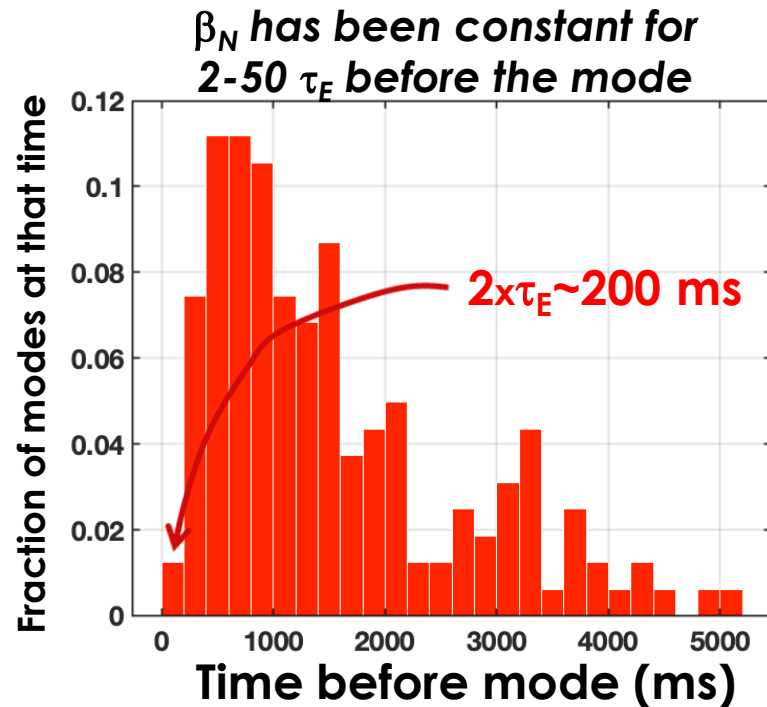
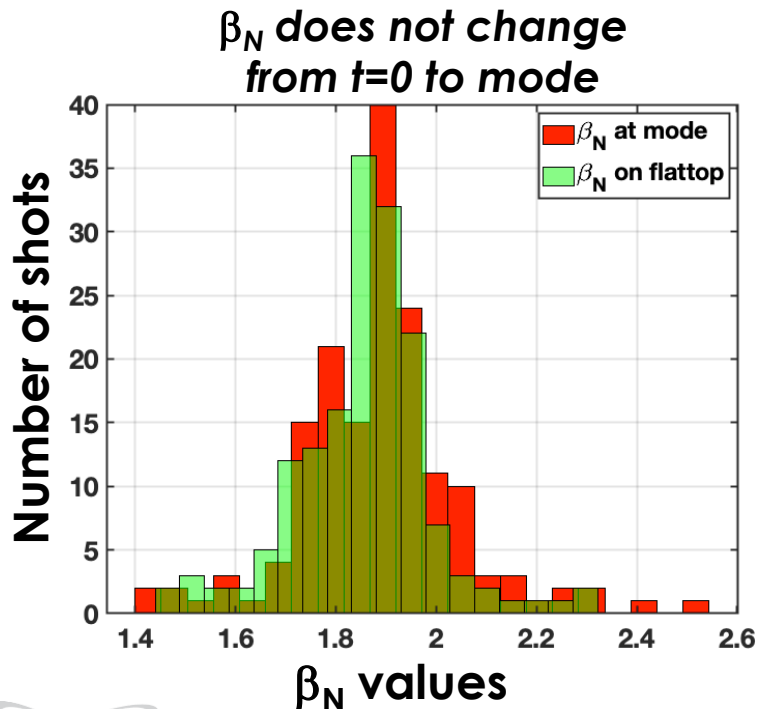
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# Pressure: The IBS Instabilities Are **Not** Due to a $\beta_N$ Limit

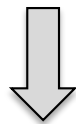
- The modes appear after  $>10 \tau_E$  at constant pressure and pressure gradient



If it they were pressure-driven, it would tear in the first 100 ms

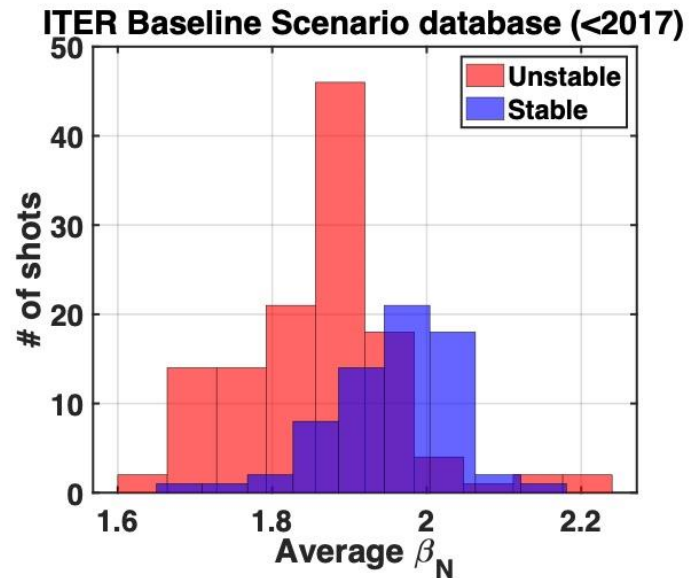
# Pressure: The IBS Instabilities Are **Not** Due to a $\beta_N$ Limit

- The modes appear after  $>10 \tau_E$  at **constant pressure and pressure gradient**
- The  $\beta_N$  is **low**  $\rightarrow \beta_N \ll \beta_{N, \text{no-wall}} \rightarrow$  not an ideal limit
- **Lower**  $\beta_N$  does not lead to better stability, **higher**  $\beta_N$  is not more unstable



**There is no  $\beta_N$  threshold for instability**

*The pressure is not the cause of the 2/1 modes in the IBS*



# Higher $\beta$ Shots Are Stable for $>10 \tau_E \rightarrow$ The Pressure is Relaxed, the Current is Not!

Timing of modes  $\rightarrow$  type of drive

**Still not a pressure limit:**

$\rightarrow$   $p$  higher from the start

$\rightarrow$  Mode hits  $>1$  s later!

$\rightarrow J_{ped}$  increases on  $\tau_R \sim 1$  s

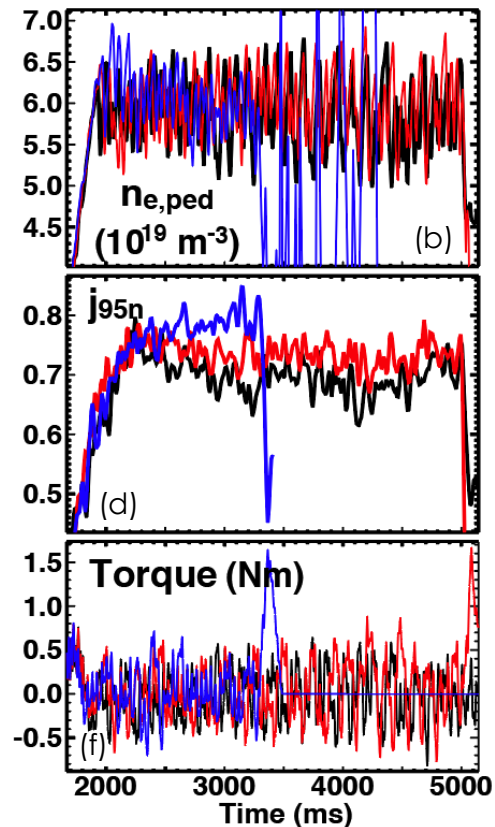
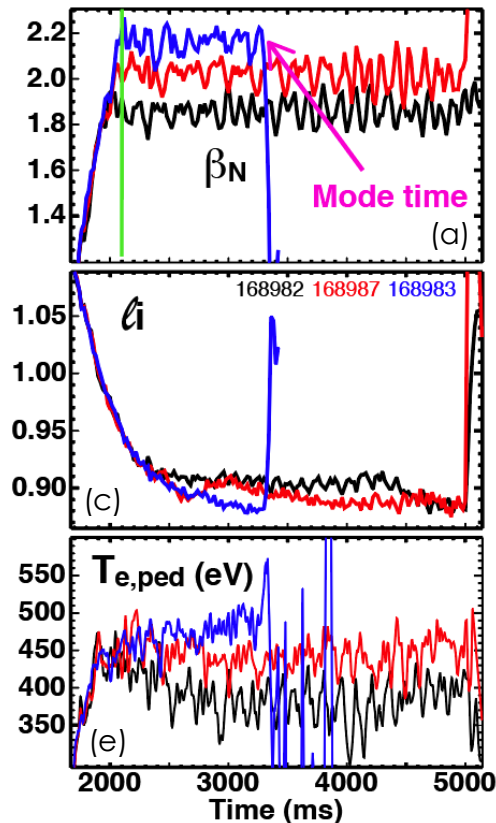
More consistent explanation:

For a given  $J$  profile:

higher  $\beta_N =$

$\rightarrow$  higher pedestal (bootstrap)

$\rightarrow$  more unstable





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*Correlation is NOT causation*

The onset condition from  $\nabla\Omega$  violates the onset time requirements  $\rightarrow$   
Eliminated (in backup slides if there is interest)

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**All inconsistent with the mode onset timing and do not represent all the database  $\rightarrow$  *Eliminated* (in backup slides if there is interest)**

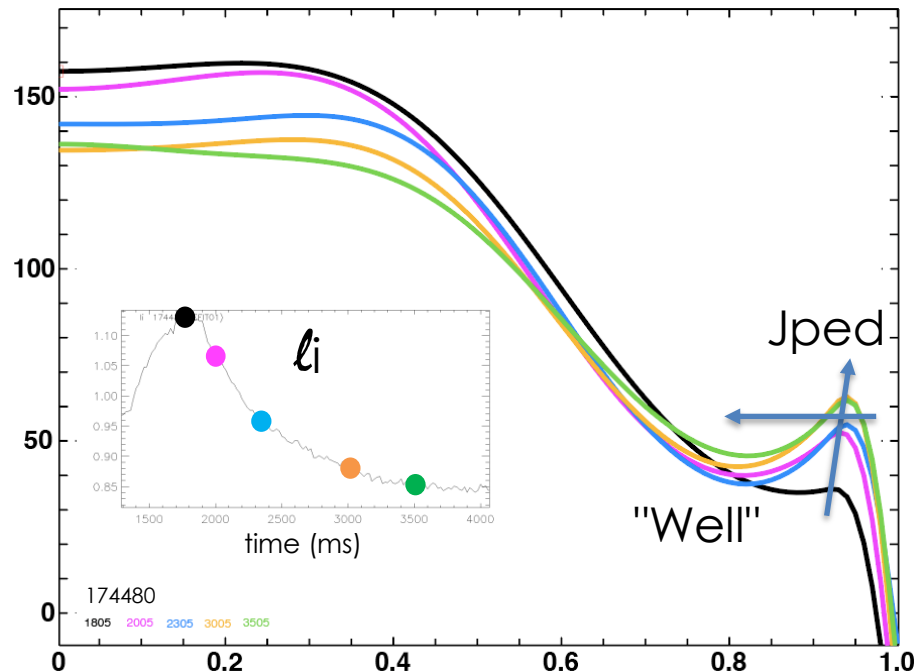
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*This actually worked...*

# Ohmic and Bootstrap Currents Create the Edge Peak and the Current "Well" That Evolves from the H-mode Transition

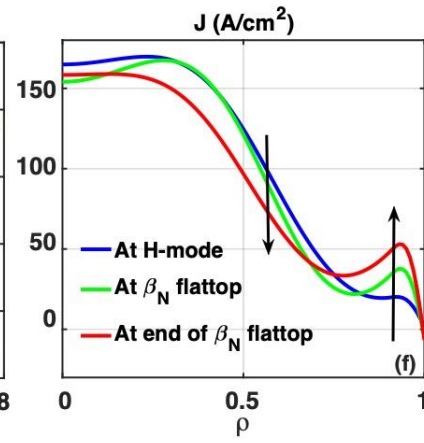
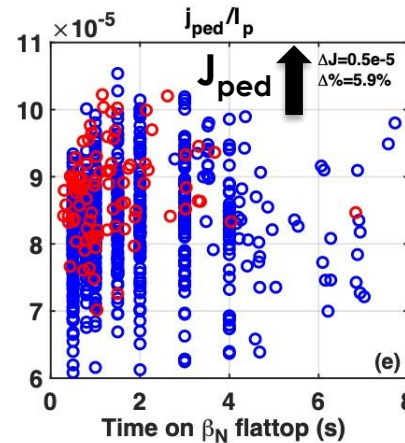
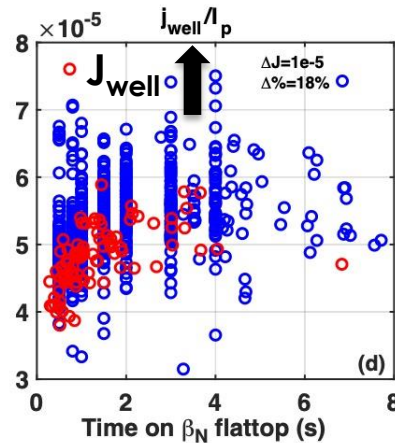
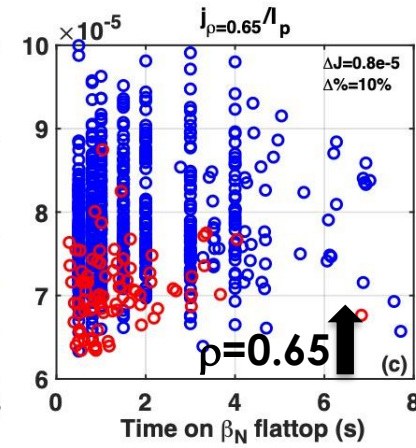
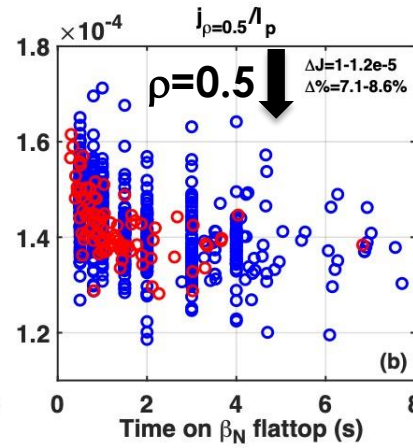
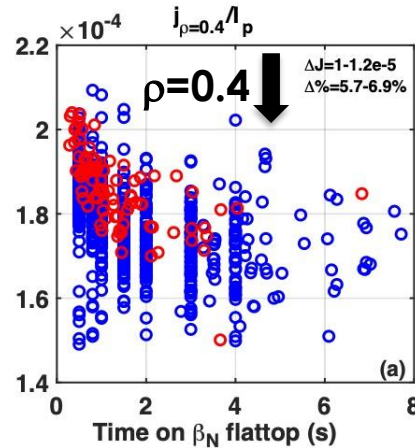
- At the H-mode transition  $J_{boot}$  is formed  
→ peak at  $\rho \sim 0.92$
- Ohmic current comes from the edge and diffuses inwards  
→ the "well" at  $\rho \sim 0.8$  slowly fills
- The core is fixed by sawteeth  $\rho < 0.45$  ( $q_{min} \sim 0.9-1$ )



# The Current Profile Grows in the Outer Region and Reduces in the Core During the $\beta_N$ flattop

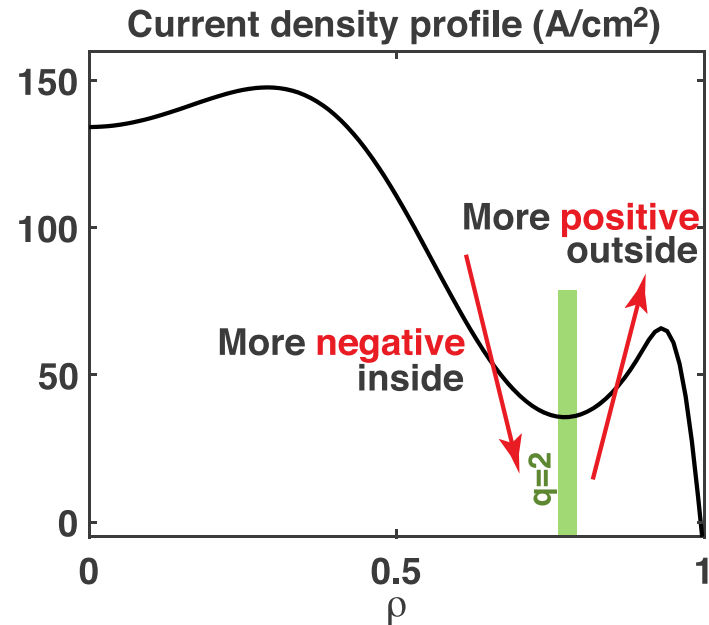
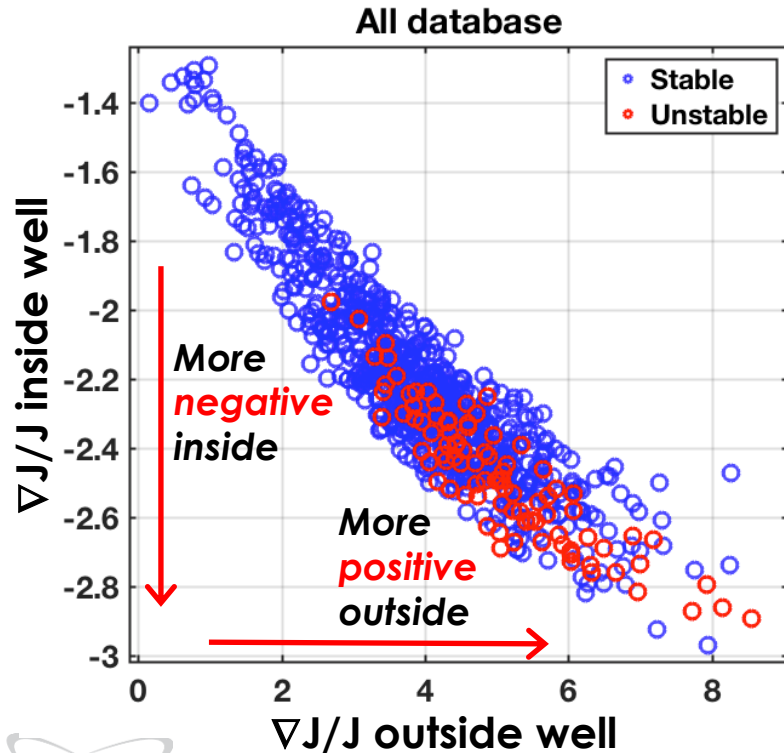
- Stable time slice
- Time of 2/1 mode

- J evolves for  $\sim 2$  s after flattop
- **Stable** and **unstable** shots have a similar evolution
- ...but the initial conditions are on average different



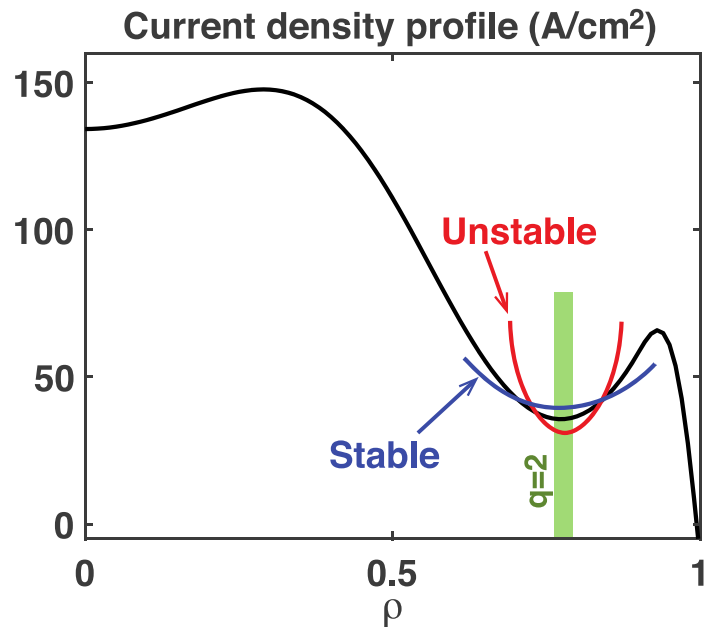
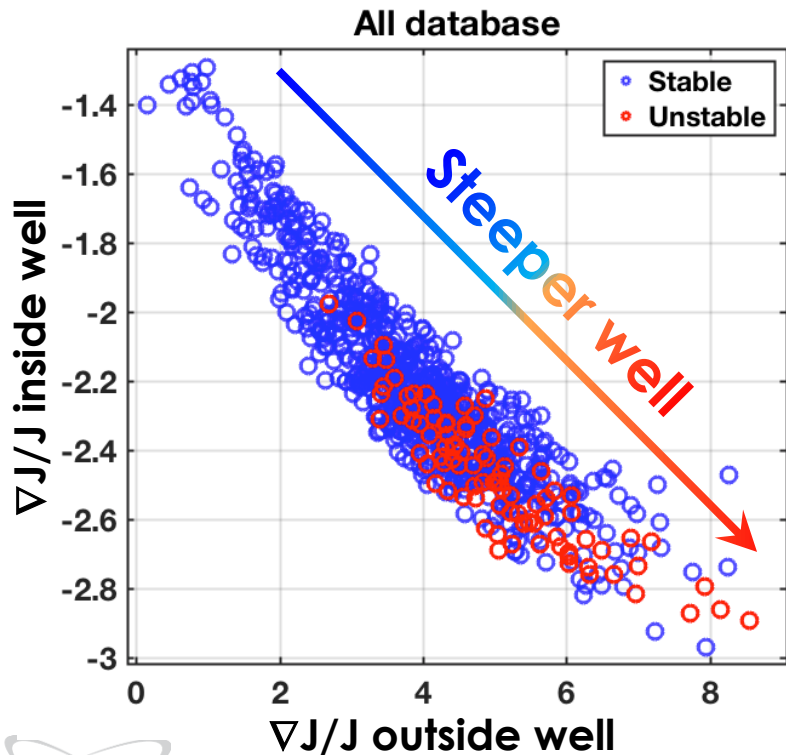
# J at the Mode Onset has a Steeper Well Around the q=2 Surface

- Both  $\nabla J_s$  are larger in magnitude at the times of the mode onset



# J at the Mode Onset has a Steeper Well Around the $q=2$ Surface

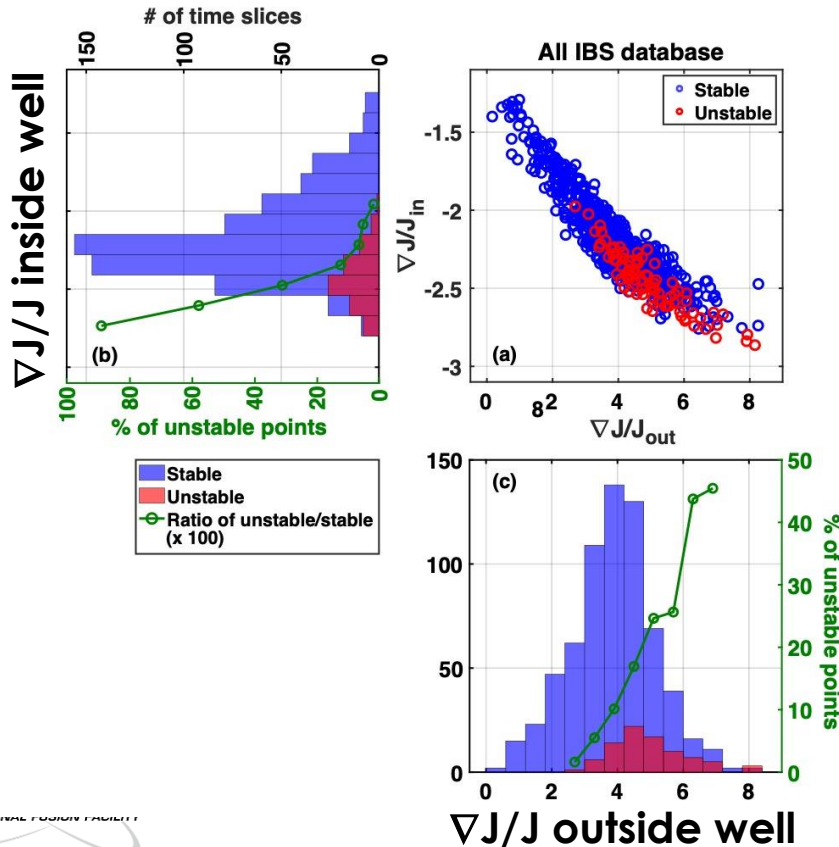
- A steeper "well" in the current profile is likely the cause of the instabilities



- Consistent with time scales  $\tau_R$
- Independent of seeding

# Separating Stable vs Unstable TIMES Shows the Correlation Between the Current Profile and the Instability

- **Unstable points** fall predominantly in the lower right region (larger gradients)

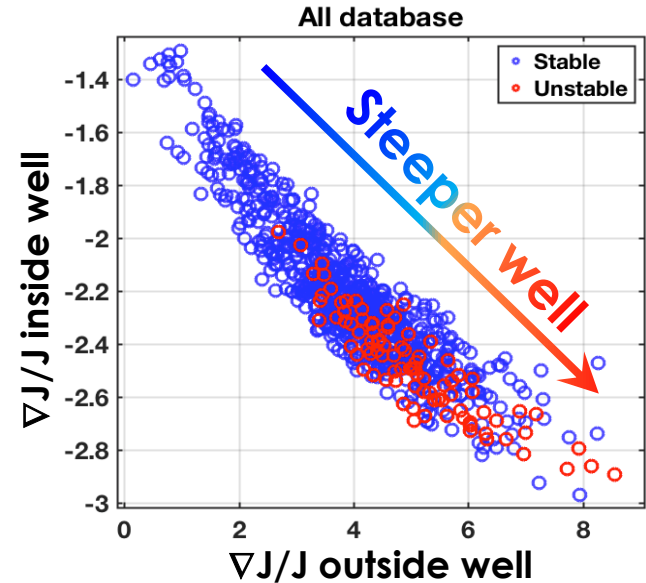


- Histograms allow to see the **whole database** of time slices ( $10^4$  hidden points)
- **Statistics** on the calculated ratios show this is meaningful (quantifiable)



# Changes in the Current Profile Affect the Classical Tearing Index $\Delta'$

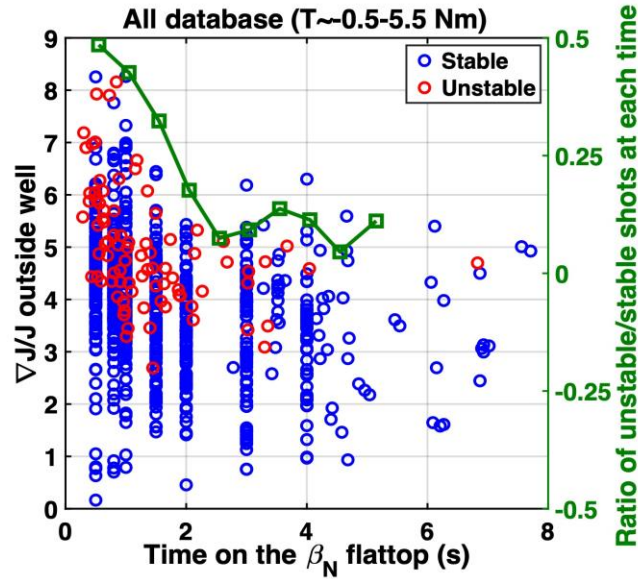
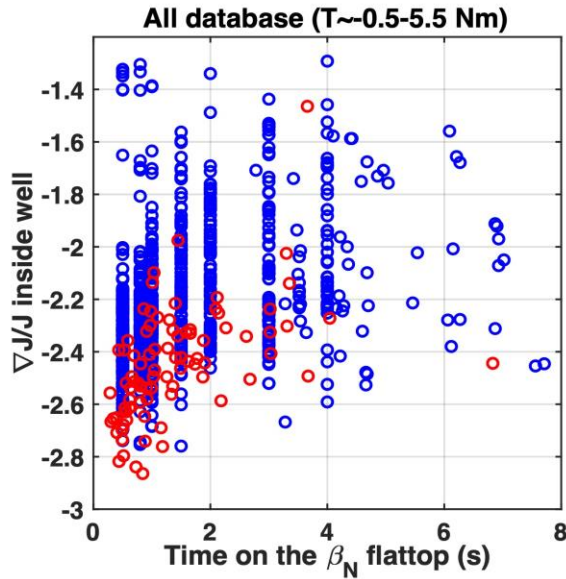
- This suggests stable and unstable times have different  $\Delta'$  : classical drive
- $\Delta'$  is a GLOBAL parameter, determined by all the current profile
- $\Delta' > 0$  is necessary, not sufficient for instability:
  - $\Delta'$  trends determine if more/less stable
  - For instability,  $\Delta' > \Delta'_{crit}$  (inner layer physics)



*The critical  $\Delta'$  for instability is likely the cause of the overlapping region of J (local  $T_e$ ?)*

# Timing of the Modes Helps Us Find a Solution: the Unstable Shots Are Separated in $\nabla J$ from the Start

- More stable current profile late, **fewer unstable shots after ~1 s**  
→ if we **solve the access problem**, high probability of remaining stable



67% of the Instabilities Occur Before 1.3 s on the  $\beta_N$  Flattop ( $1-1.5 \tau_R$ )

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**2. MHD stability and disruptions**

- Cause of the instabilities
- **Solution and new scenario**

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# Methodology to Stabilize the IBS Modes Illustrates Actuators That Can be Used in All Scenarios to Change the Current Profile

Changes  $l_i$ , current penetration



Changes ELM frequency, pedestal (not density)

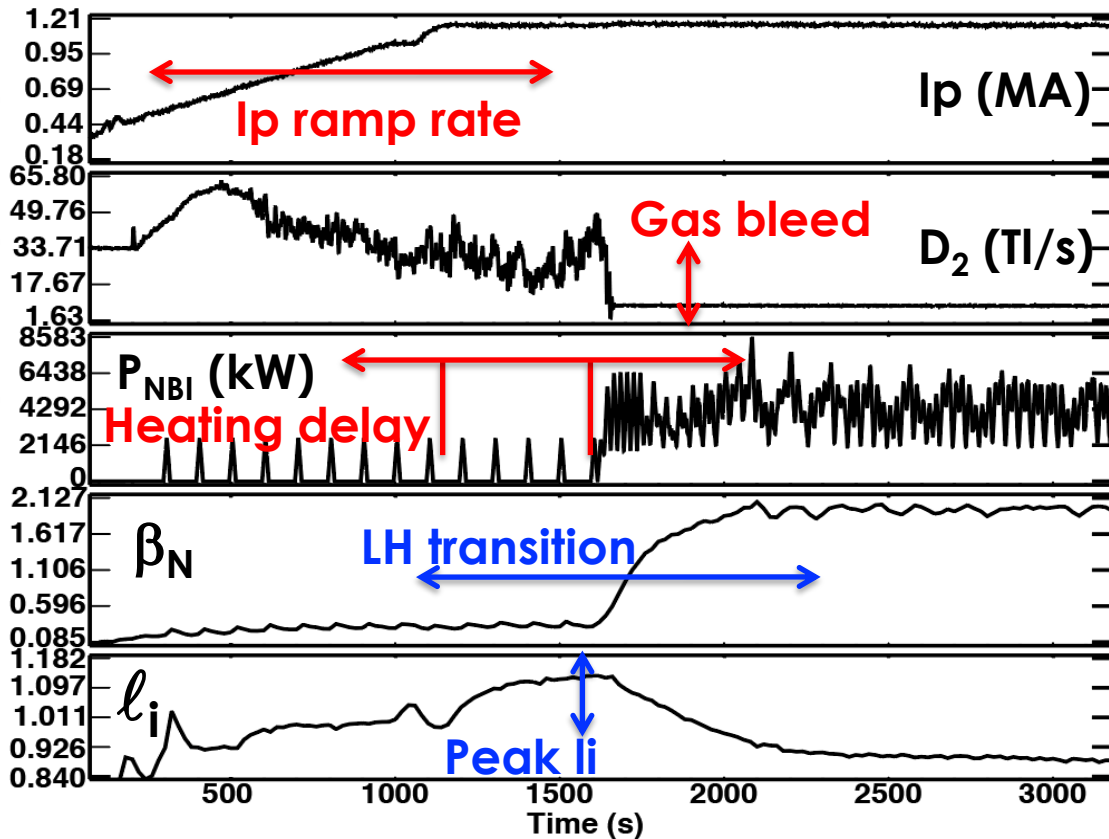


Changes  $l_i$ , current penetration



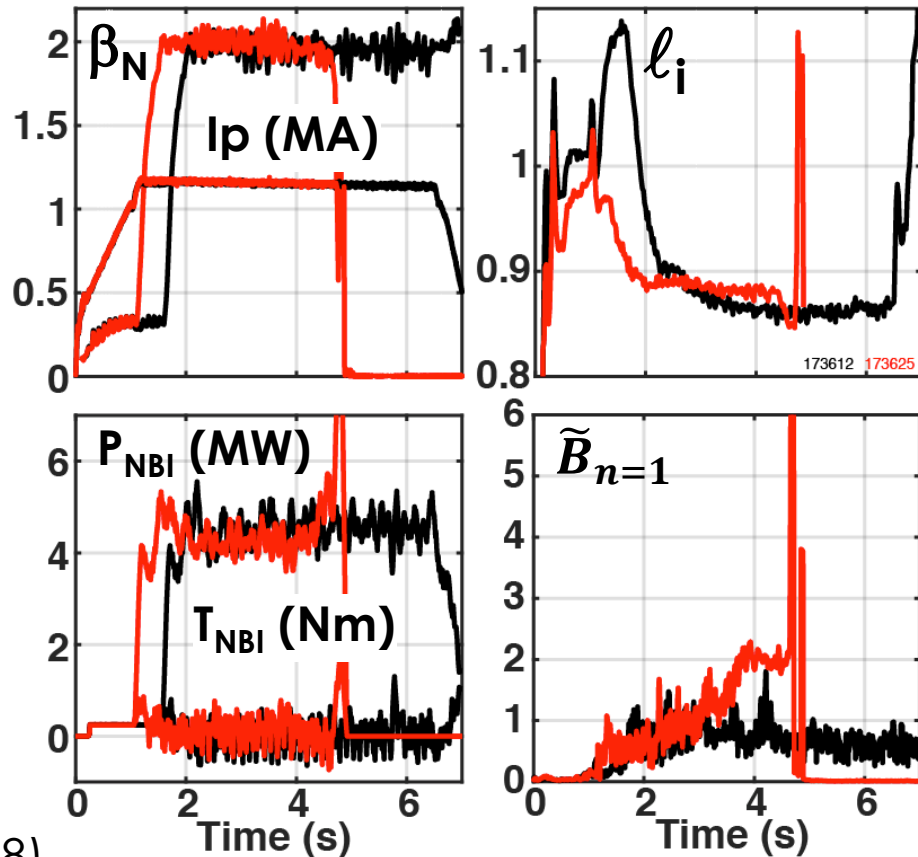
Allow  $J$  time to penetrate into the core  $\rightarrow$  starts with lower pedestal

Fixed zero torque



# Applied Modifications to Show Causality. (1) Heating delay

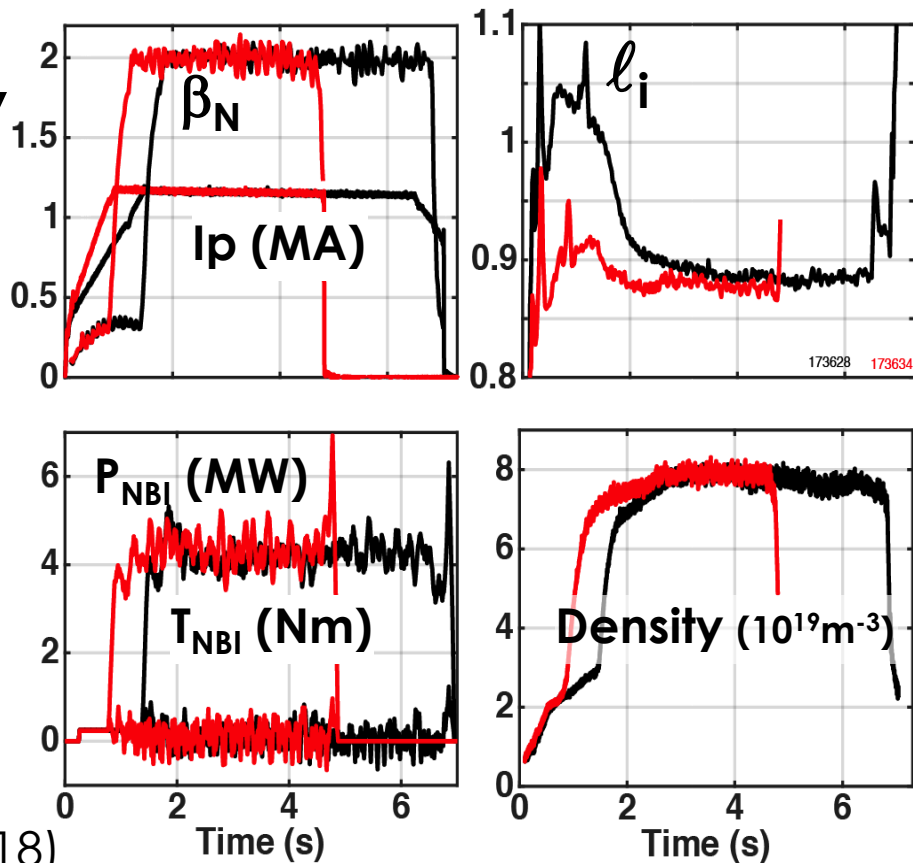
- Database of pulses with only change to the H-mode transition time shows the late timing is robustly stable
- Trajectory of  $l_i$  shows current profile evolution is different
  - $l_i$  is not sufficient to predict stability



(T.C. Luce 2018)

# Applied Modifications to Show Causality. (2) $I_p$ ramp rate

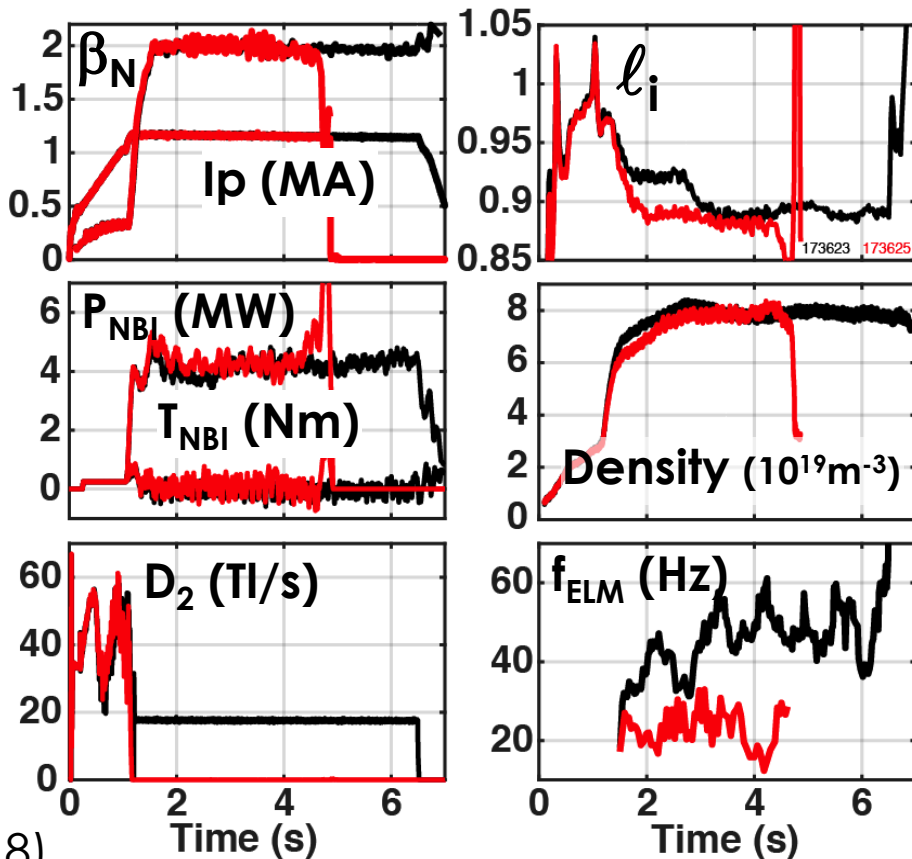
- Slower  $I_p$  ramp rates are robustly stable – similar effect as heating delay
- Combination of  $I_p$  ramp and heating time changes can tailor the stability to the hardware requirements



(T.C. Luce 2018)

# Applied Modifications to Show Causality. (3) D<sub>2</sub> gas "bleed"

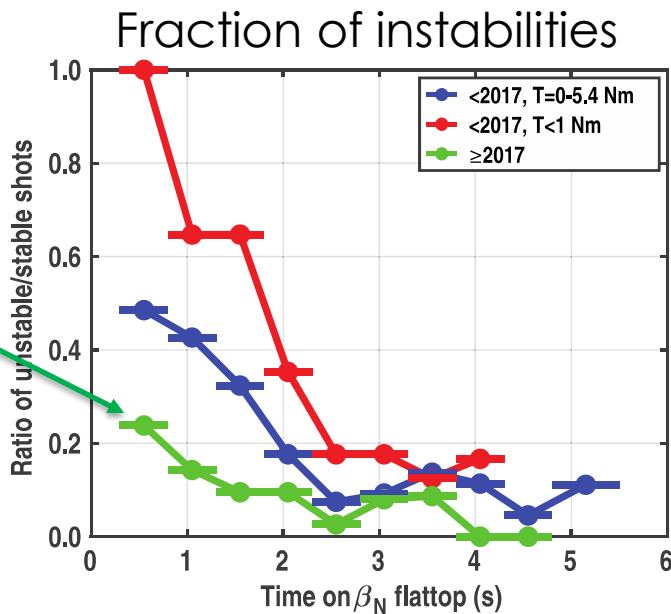
- Modest gas "bleed" eliminates LATE modes
- Results in more regular and more frequent ELMs
- Smaller ELMs = lower pedestal!
- Little difference in density



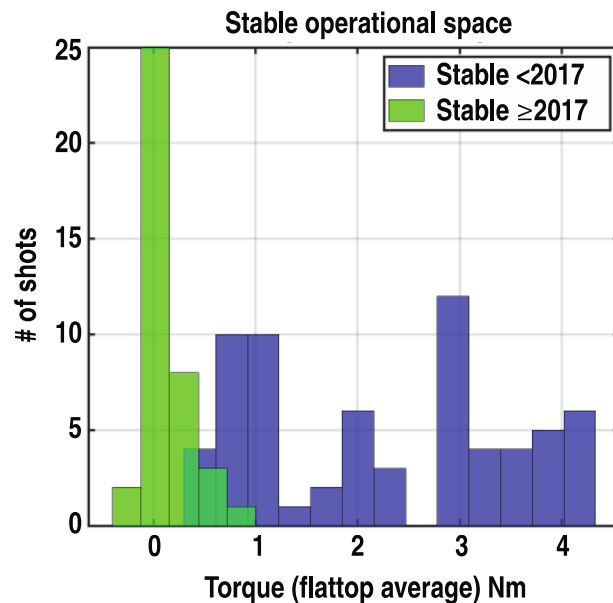
(T.C. Luce 2018)

# The Recipe is Robust and Repeatable Under a Variety of Conditions

- Different  $I_p$ ,  $B_T$ ,  $n_e$ , gas, impurities (Kr, Xe, C6, W)
- Heating mix (ECH, NBI)
- Open/closed divertor (USN, LSN)
- It did not go to a different density,  $\Omega$ ,  $\Delta\Omega$  regime



This is NOT the intrinsic instability!  
(intentional disruptions)





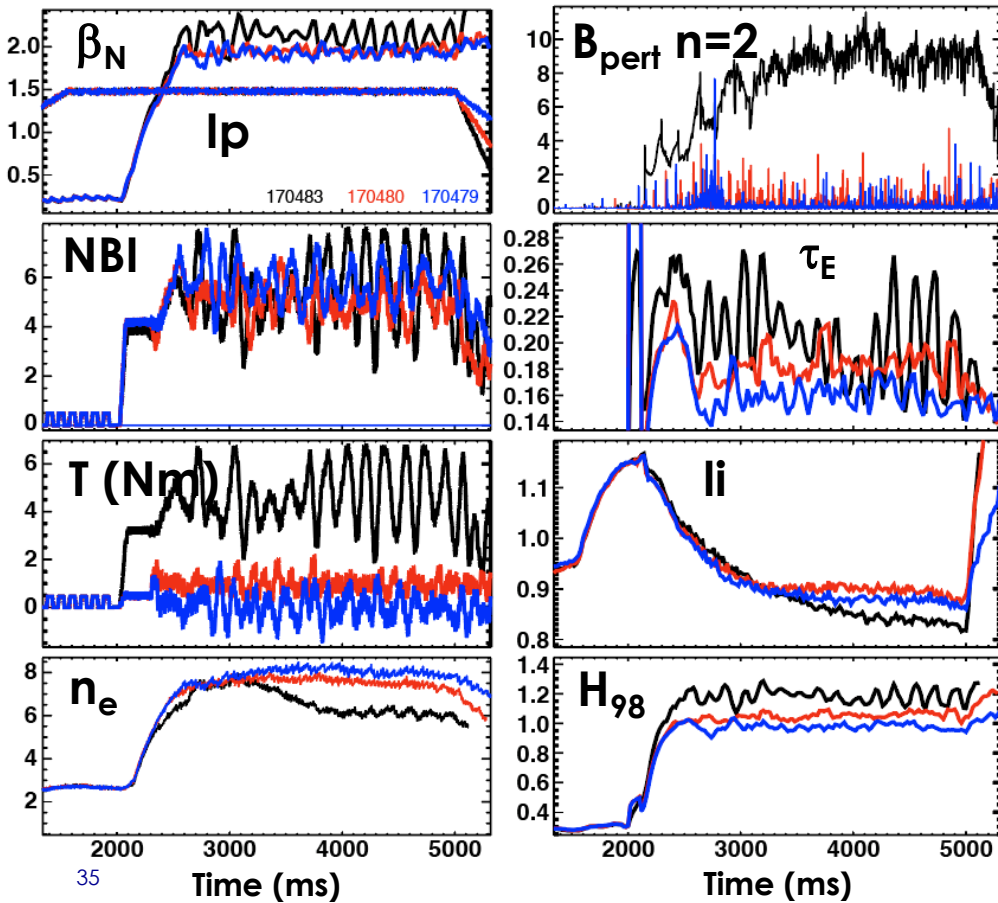
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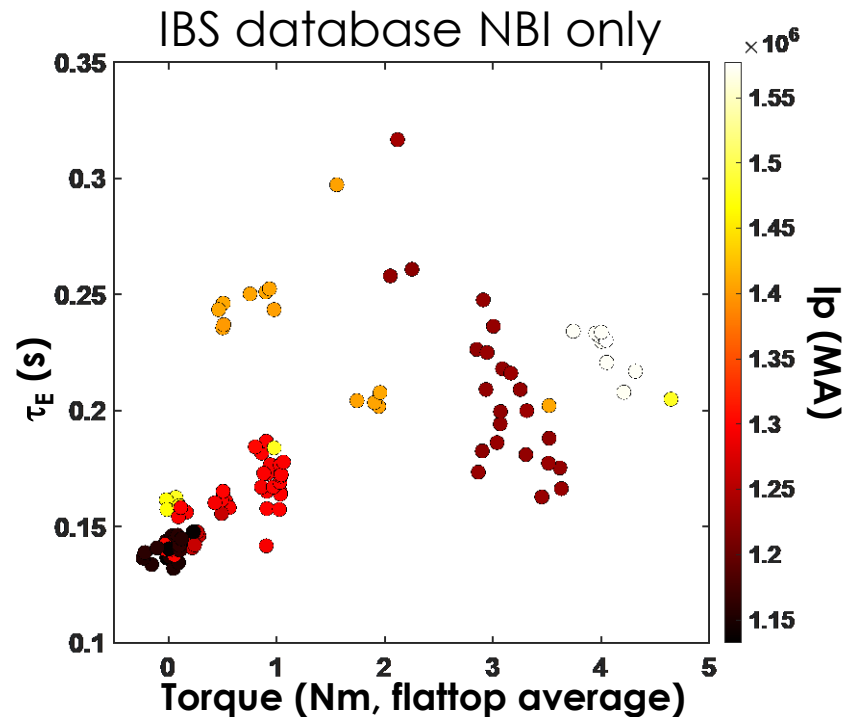
# Metrics for Evaluation of the Performance

- **Fusion power in DT plasmas of interest in ITER will have:**  
 $P_{\text{fus}} \propto \langle p^2 \rangle \propto \beta_T^2 (\%)$  at fixed B (makes  $P_{\text{fus}}$  dimensionless)
  - In ITER,  $P_{\text{fus}} = 500$  MW at  $B=5.3$  T requires  $\beta_T=2.55\%$
- $Q_{\text{fus}} \propto \langle nT \rangle \tau \Rightarrow$  **use  $\beta_T \tau$  as a proxy for gain (not dimensionless)**
  - Can also use  $G \equiv \beta_N H_{89} / q_{95}^2$  as a proxy for gain, but the accuracy of a confinement scaling is assumed
  - ITER  $Q=10$  requires  **$G=0.38-0.42$**  (depends on precise value assumed for  $q_{95}$  at 15 MA)
- **Will also show the standard stability and confinement metrics ( $\beta_N, H_{98y2}$ )**

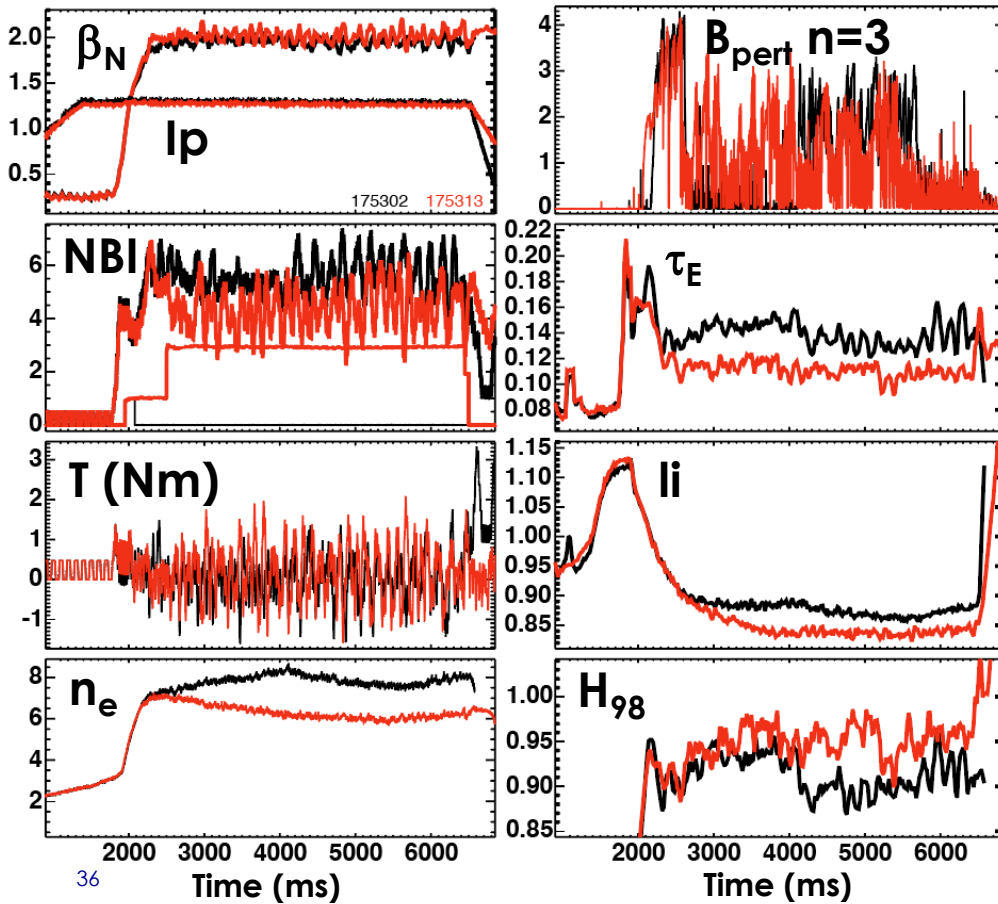
# Confinement Gets Worse at Lower Torque



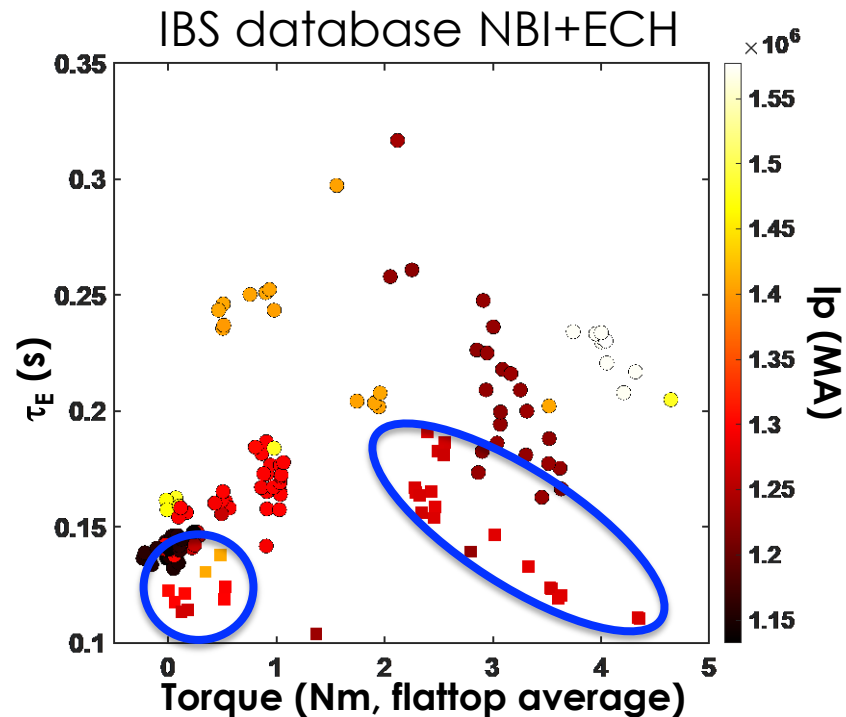
- But other factors apply ( $I_p$ , density, etc...)



# ECH in Dominant NBI Heated Plasmas Reduces Confinement



- ECH degrades both heat and particle confinement



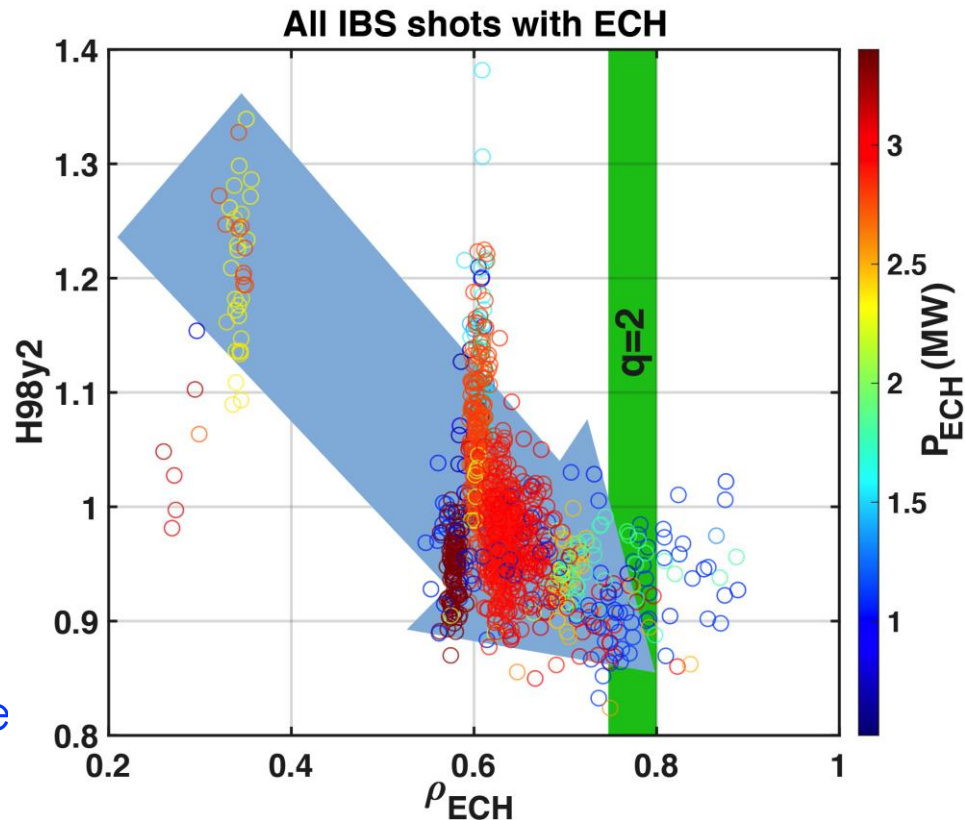
# The ECH Location Also Strongly Affects the Confinement

Heating efficiency drops dramatically with off-axis heating

$$\eta_H \sim 1 - \rho^2$$

**Regardless of  $I_p$ , torque:  $\rho_{ECH} > 0.75$**

- Reduces confinement below standard H-mode levels  
*(Not a power degradation effect)*
- Direct ECCD stabilization is not compatible with Q=10 performance



# Drop in Confinement With Off-Axis Heating Can Be Seen Also Dynamically With RT Mirror Steering of the Gyrotrons

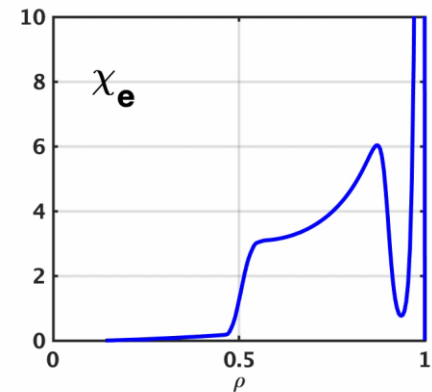
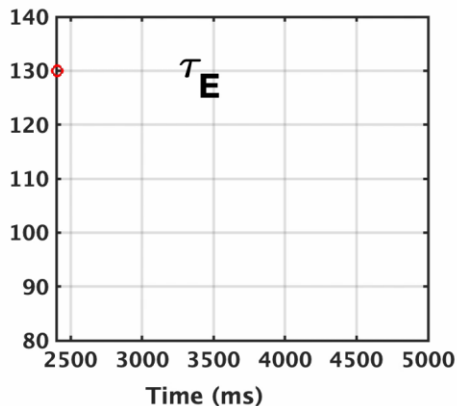
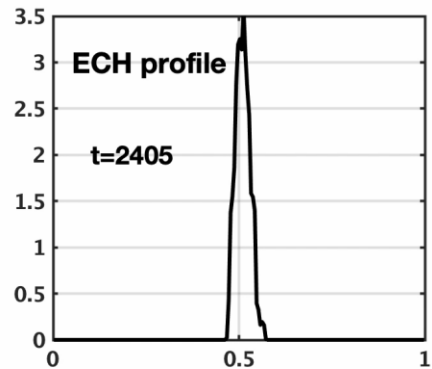
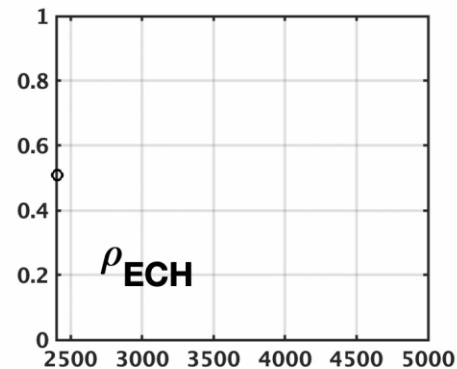
- Moving 3 MW of ECH from  $\rho=0.5$  to  $\rho=0.8$  decreases  $\tau_E$  by 25-30%,  $H_{98y2}$  by 15-18%

$$\eta_{heating} \approx 1 - \rho_{ECH}^2$$

Expect 50% drop in  $\tau_E$ , observe 25%



Loss of heating efficiency  
compensated by transport improvement



# Drop in Confinement With Off-Axis Heating Can Be Seen Also Dynamically With RT Mirror Steering of the Gyrotrons

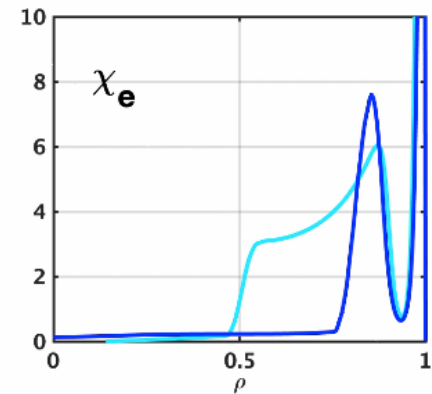
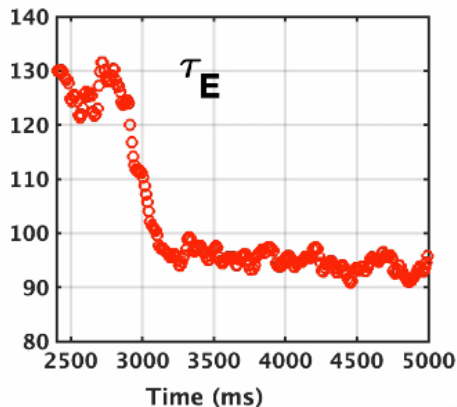
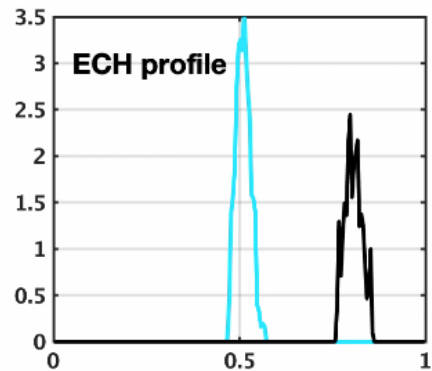
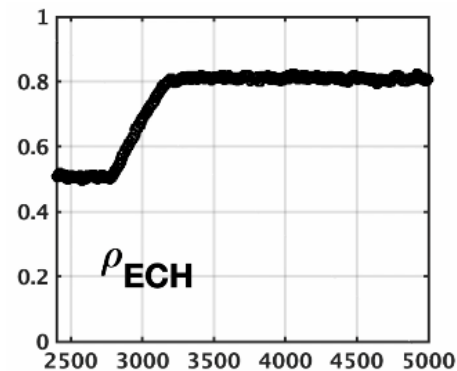
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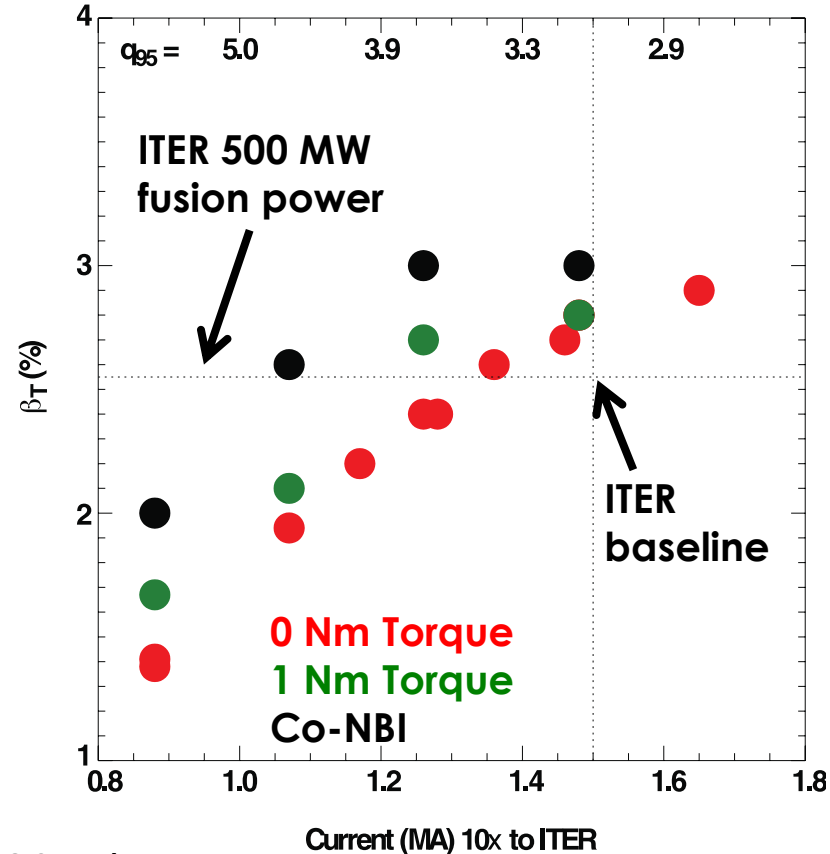
# What we discuss today

1. Standard IBS pulse design
2. MHD stability and disruptions
  - Cause of the instabilities
  - Solution and new scenario
3. Confinement trends
- 4. Lower current and heated access options**



# Performance to Reach 500 MW of Fusion Power in ITER Achieved at All Torque Levels and Lower Current

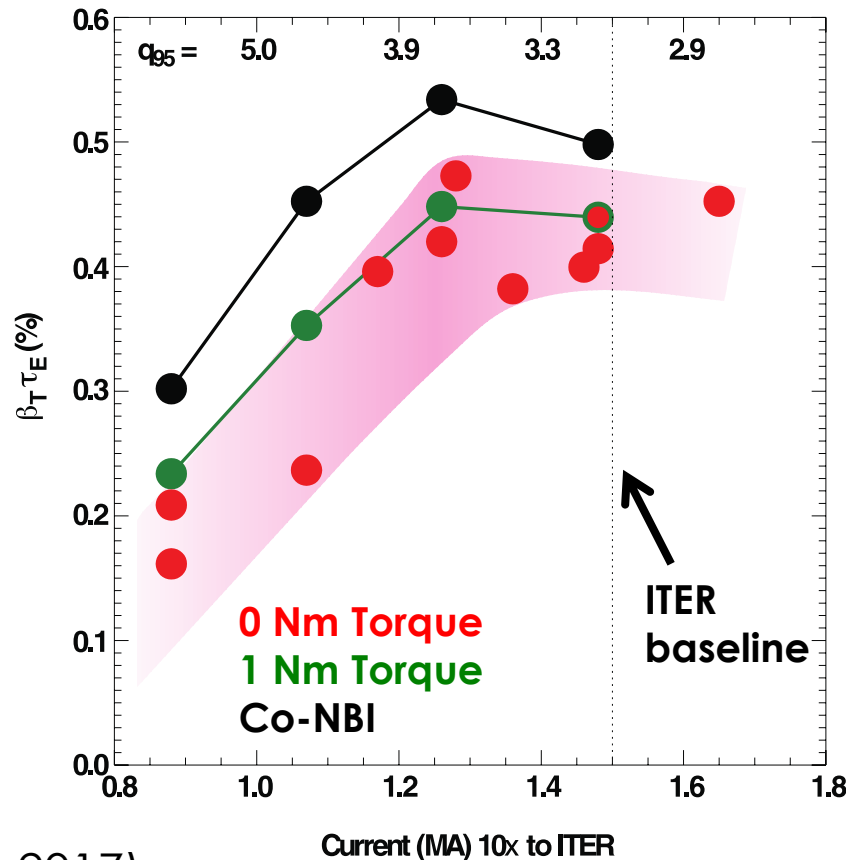
- With co-NBI, the goal is reached by 11 MA equivalent
- With 0 Nm torque, 13.5 MA may be sufficient
- For co-NBI, the achieved  $\beta$  does not increase above 12.5 MA



(T.C. Luce 2017)

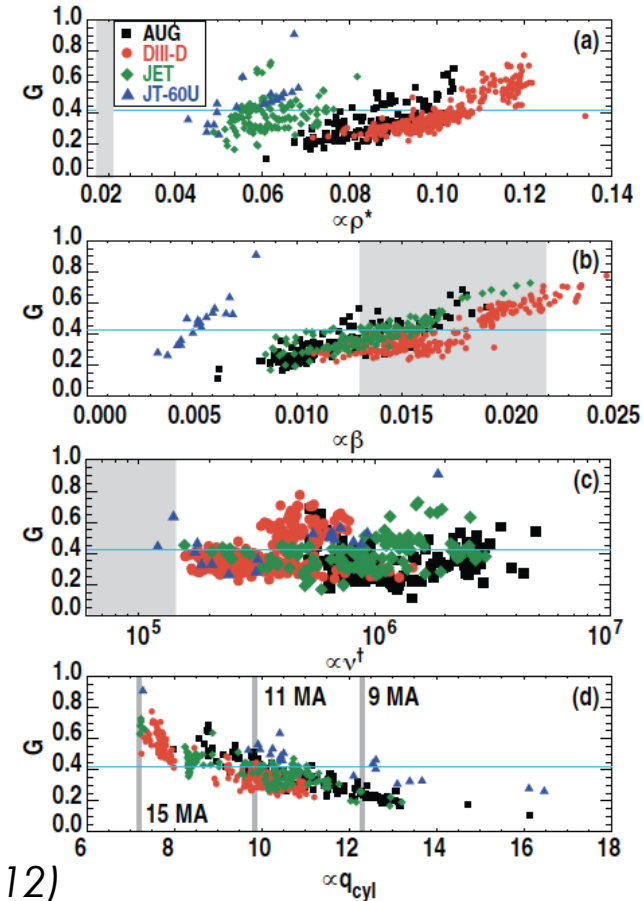
# Gain Metric ( $\beta\tau$ ) Does Not Improve Above 13 MA Equivalent Current

- **Curves at all torque levels have similar shapes**
  - Effect is not likely due to ExB shear
- **Increase in gain seems to saturate around 13 MA**
  - Corresponds to  $q_{95} \approx 3.7$
  - Previously seen on DIII-D, but not explained [Schissel, et al., NF 32, 107 (1992)]



# The Advanced Inductive Scenario Constitutes a Promising Alternative for ITER's Q=10 Mission at Lower Current

- **Advanced Inductive** (A.I.) scenario has demonstrated good performance projecting to Q=10:
  - $H_{98y2} \sim 1.2-1.5$ ,  $\beta_N = 2.4-3.0 \rightarrow G \sim 0.4$
  - At  $q_{95} \sim 3.8-4.8$  and  $T_i/T_e \sim 1-1.6$
  - In DIII-D, AUG, JT-60U and JET
- This scenario has a different access to H-mode: heated ramp
- Needs to be extended to SN shape and low T

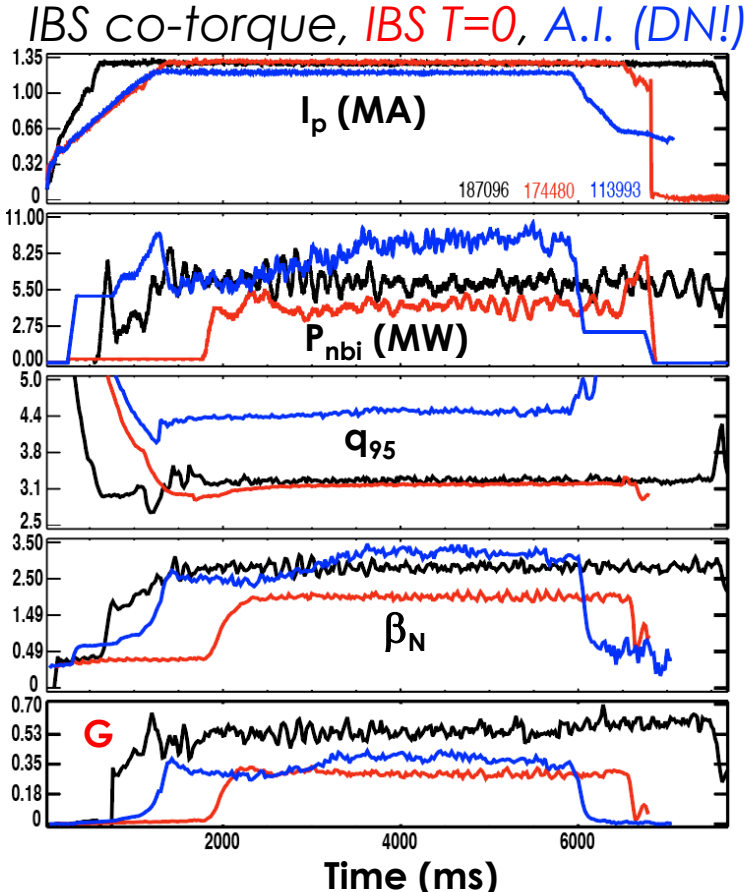
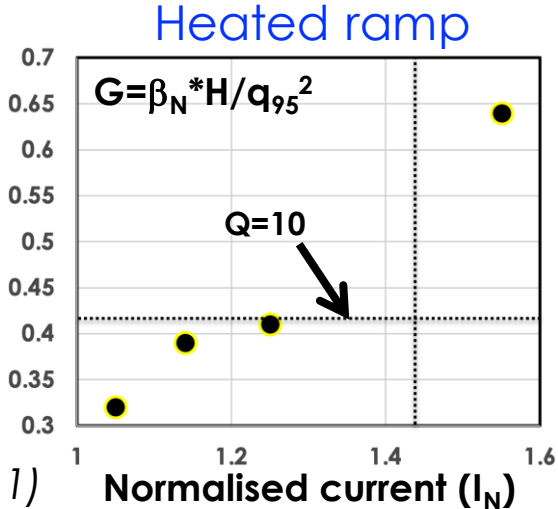


(Luce 2012)

# Ohmic vs A.I. Heated Access Has Significant Impact on the Stability and Performance of the Burn Phase

- Heating in the Ip ramp usually gives access to higher stable flat-top  $\beta_N$
- A.I. scenario can operate at reduced  $q_{95} \sim 4.5$  because of higher MHD limits
- G is maintained thanks to higher  $\beta_N$  and high  $H_{98}$

JET hybrid scenario results in next talk!



# What we have not covered (a lot)

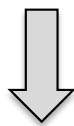
- Tungsten and the impact of intrinsic radiation
- ITER-like controls: slower, fewer?
- Projections (no time today...)
- Differences with ITER parameters:
  - Higher temperature ( $I_p$ , power)
  - Lower collisionality
  - Rotation is relatively unknown → confinement?
  - Core-edge integration (divertor impurities and detachment)
  - RMP ELM suppression
  - Fuelling: less core, more pellets?



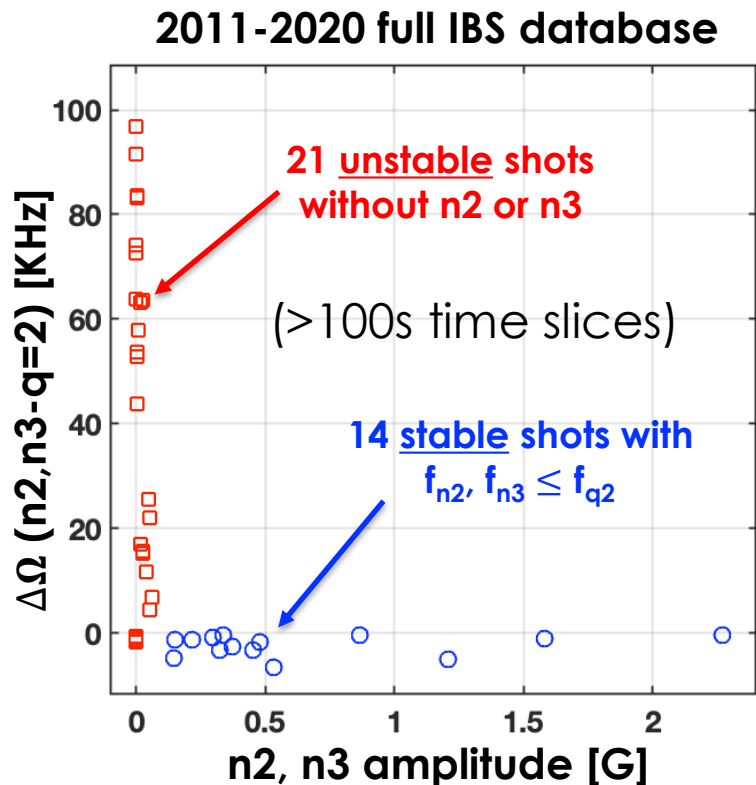
# Stability: seeds (not)

# Experience, and DB analysis, show that the presence of other $n > 1$ modes does not separate stable/unstable plasmas

- Many of the 2/1 modes are **not coincident** with any other  $n > 1$  mode
- Many of the stable shots have zero rotation differential between  $n_2, n_3$  islands and the  $q=2$  surface - but **no 2/1 mode**



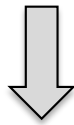
Higher  $n$  modes cannot be a general cause of the instabilities in the IBS



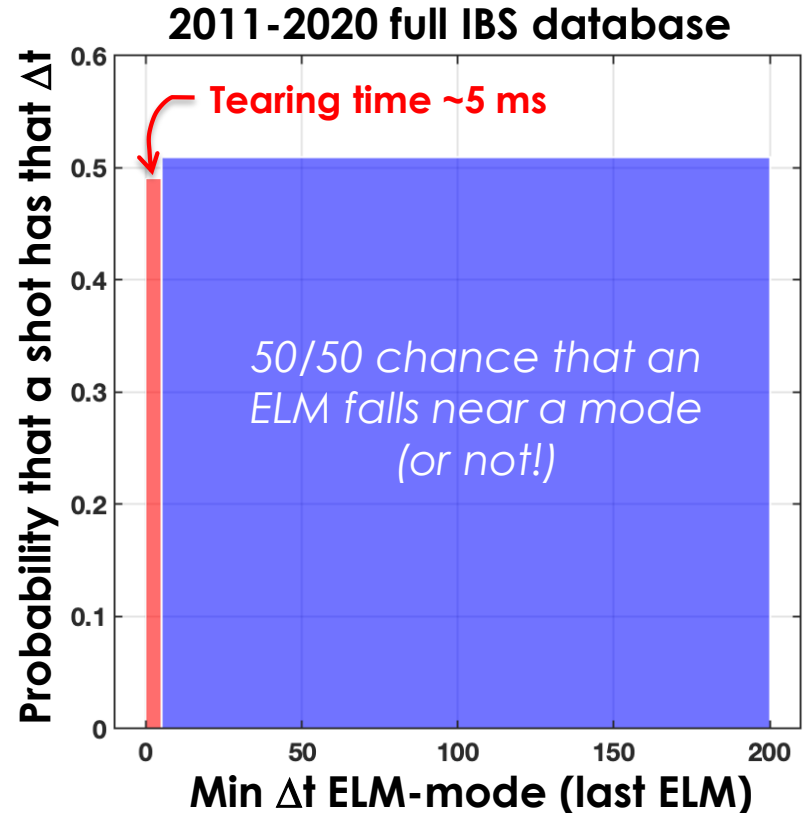


# Millions of seeds... and very few modes $\rightarrow$ seeding is not the main cause of the 2/1 instabilities

- There are 18054 ELMs in 273 shots – only 162 modes  $\square$
- Stable shots have the same ELMs and sawteeth as the shots with a mode
- There are hundreds of ELMs and tens of sawteeth before a mode



**Focussing on the last lone ELM or sawtooth will not solve the problem**



# Stability: rotation (not)

# J, p represent the equilibrium, $\Omega$ may affect mode coupling or locking

## Sources of free energy

- **Pressure, current** (and their profile shapes)
- $\Delta'$  is a global parameter **function of p and J** = free energy for tearing  
 $\Delta' > 0$  *Necessary but not sufficient:*
- Instability threshold  $\Delta'_{crit} > 0$  in toroidal geometry: function of inner layer physics →  $T_e, n_e, \nabla T_e$ , etc

## Rotation and its gradient:

- $\Omega$  **not a source of free energy**
- **at rational surface**, large  $\Omega$  *gradient* destabilizing → need very large island (not our case)
- **between rational surfaces**, large  $\Omega$  *gradient* believed to decouple equal n surfaces → only 1/1 and 3/1 surfaces can couple

# Unlike $J$ , $\Omega$ and $\Delta\Omega$ Evolve Fast, or Are Fixed at Low Torque

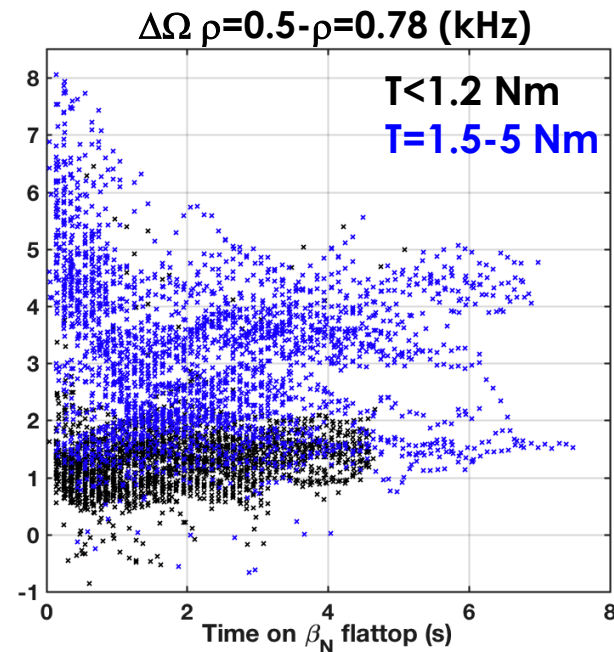
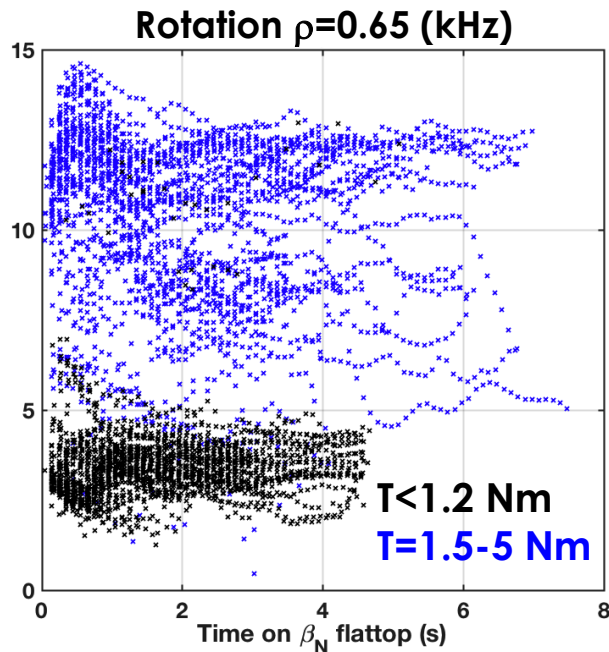
- Time-of-mode plot shows that stability gets better in time (fewer modes late)

$\Omega$  and  $\Delta\Omega$  either decrease (worse) or stay the same

The equilibrium becomes more stable,  $\Omega$  does not

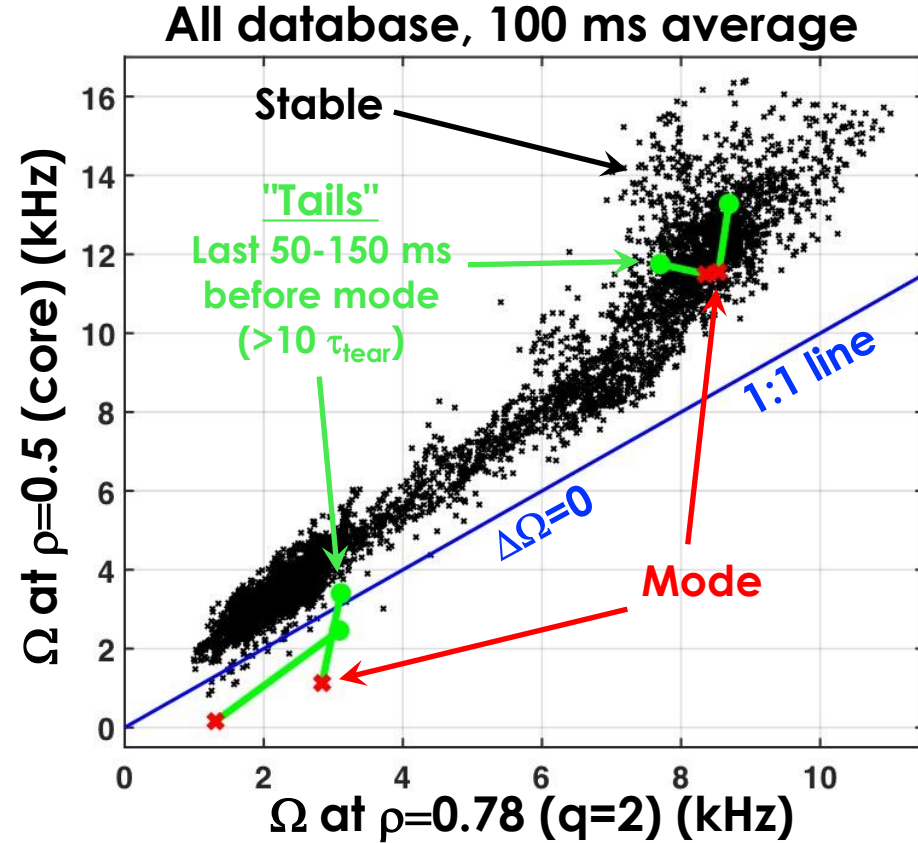


**Hypothesis of destabilization by  $\Delta\Omega$  inconsistent with sparser modes late**



# Time Scales of $\Delta\Omega$ Evolution to Instability Are Too Long to Explain the Mode Onset by Surface Coupling

- Differential rotation between core and edge appears **lower** for many unstable times
- **"Tails"** show that  $\Delta\Omega$  is low for  $t > 10-30 \tau_{\text{tearing}}!$  ( $\tau_{\text{tear}} \sim 5 \text{ ms}$ )



# Lack of $\Delta\Omega$ is Not Likely To Be the Cause of the Instabilities

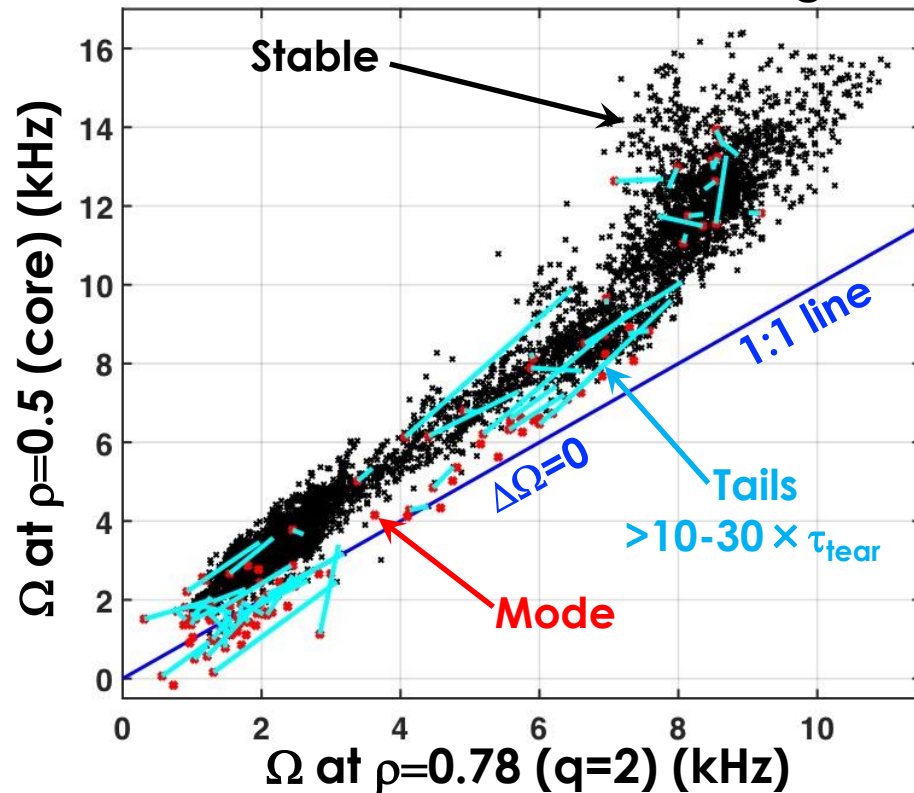
50% of the modes have

- Large  $\Delta\Omega$  (in "stable" region)
- Constant  $\Delta\Omega$  ("tail" // to 1:1)
- Low/zero  $\Delta\Omega$  for  $>10-30 \tau_{\text{tear}}$



**Mode "triggering" by lack of  $\Delta\Omega$  is inconsistent with the onset time scale**

All database, 100 ms average



# Stability: rotation is coupled to current profile

# In the Pedestal, the Time Scales are Shorter

Current → Energy, Rotation → Tearing:

- **Global resistive diffusion time**  $\tau_R \sim 800\text{-}1200$  ms,  $\tau_R \sim \Delta\rho^2$   
→ Local J changes at  $\rho \sim 0.70\text{-}0.95$  can be on  $\Delta\tau_R \sim 100$  ms
- **Energy confinement time (pressure, rotation)**  $\tau_E \sim 100\text{-}200$  ms  
→ if torque is stepped up,  $\Omega$  and  $J_{\text{pedestal}}$  change in  $\sim 200$  ms
- **Tearing time**  $\tau_T \sim 5$  ms

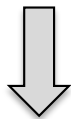
*Rotation and current profile can be coupled locally*



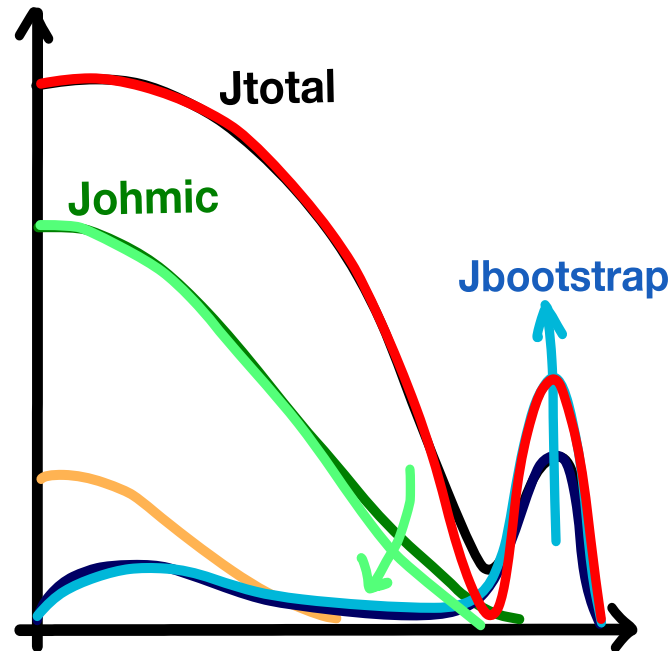
# Fixed Ip Forces the Correlation Between "Well" and Pedestal

At lower rotation/rotation shear:

- The pedestal increases (transport)
- More  $J_{boot}$
- The Ohmic drive has to reduce (fixed Ip)



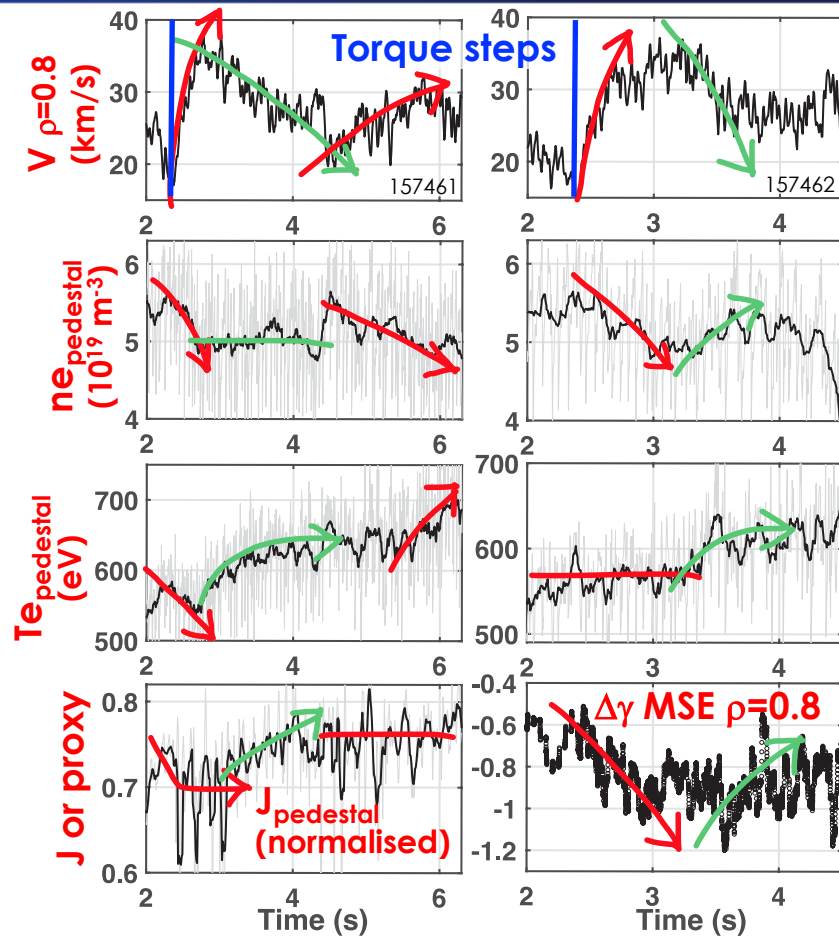
**Steeper "well" around  $q=2$**



# When the Rotation Changes, the Current Profile Changes too!

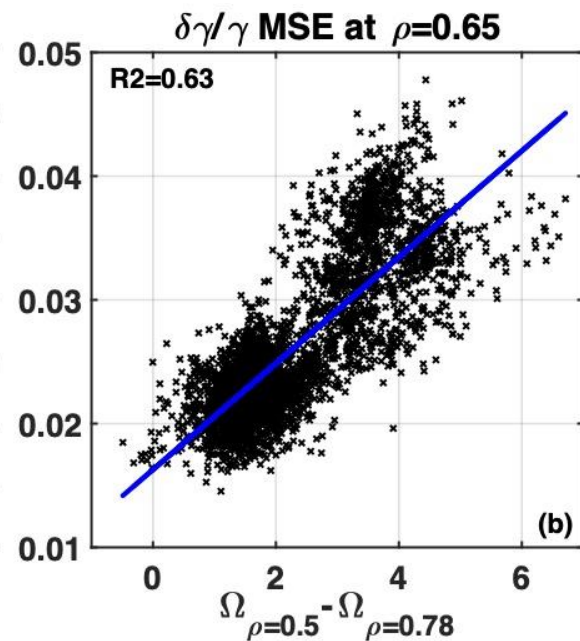
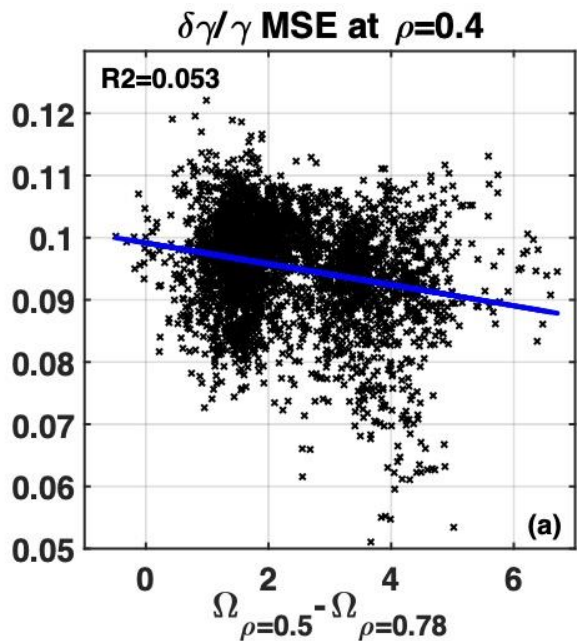
- Rotation keeps evolving after the **torque step**!
- Rotation and pedestal J are anti-correlated

Rotation	↗
$n_e$ , pedestal	↘
MSE J at "well"	↗
$J_{\text{pedestal}}$	↘



# Current and Rotation Are Also Correlated in the Global Database – raw data

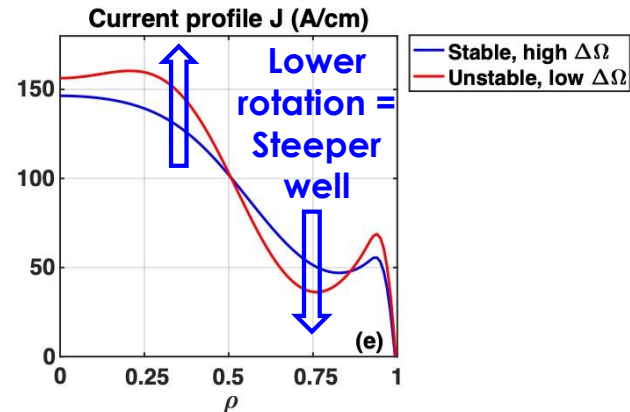
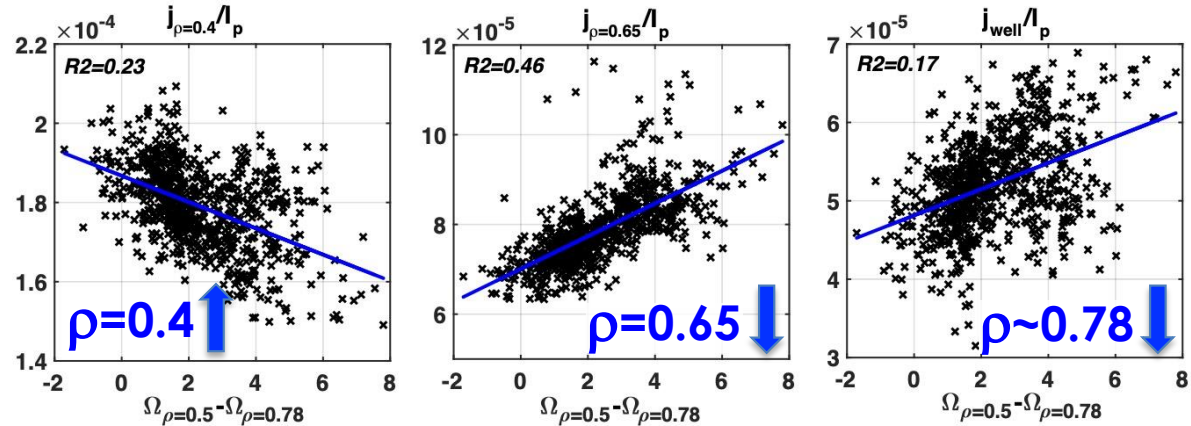
- If you don't see an effect in the raw data, it does not exist
- Raw MSE data show that  $\Omega$  and  $\Delta\Omega$  are correlated in the IBS database



# Current and Rotation Are Also Correlated in the Global Database – raw data and J (efit) reconstructions

- **MSE constrains the core, magnetics the pedestal**

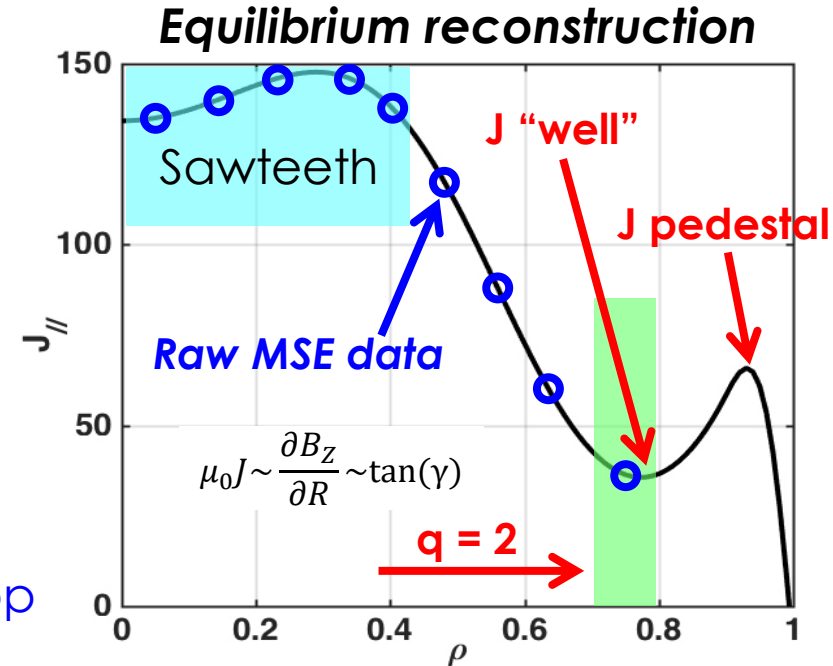
- **The total plasma current  $I_p$  is fixed  $\rightarrow$  lower J in the "well" requires higher J in the core**



# Stability: more on current profile

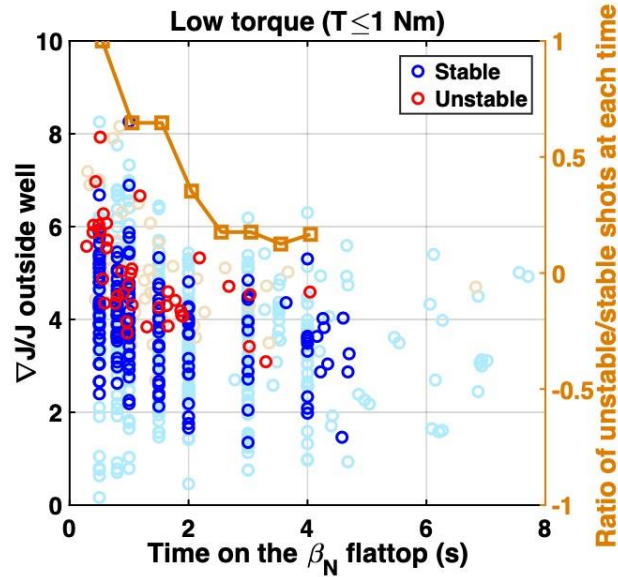
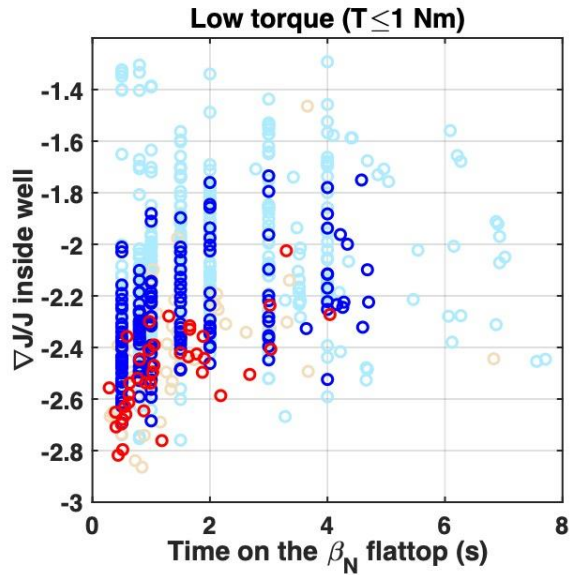
# The Current Profile of the ~330 IBS Shots can be Reconstructed and Correlated with Stability and Global Quantities

- Local current density measured by raw MSE data
- “Enhanced” fits for  $q$  and  $J$ , with MSE+magnetics and a pedestal
  - Magnetics + edge constraints describe the pedestal (similar to kinetic fit)
  - MSE constrains the core up to  $\rho \sim 0.8$
- Unstable = at time of 2/1 mode onset
- Stable = stable time slices on the  $\beta_N$  flattop (stable + unstable before mode!)



# Separation and Its Evolution Are Independent from the Applied Torque

- Low torque shots tend to be more unstable
- They start, and stay, **predominantly in the higher  $\nabla J$  regions**



$\nabla J$  separation is also independent from sawteeth, ELMs, higher m/n modes

# Represent the Evolution with Gradients Inside and Outside the J "well"

- It takes 1.5-2 s for the current profile to reach equilibrium
- Both  $J_{ped}$  and  $J_{well}$  grow, but they evolve **toward a shallower "well"**

